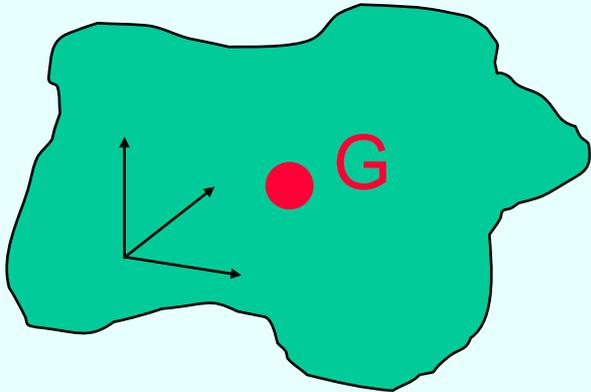


STABILITY AND TRIM OF MARINE VESSELS

*Acknowledgements to
Lt. Greg Mitchell for Slides 15-37*

Concept of Mass Center for a Rigid Body



Centroid – the point about which moments due to gravity are zero:

$$\Sigma g m_i (x_g - x_i) = 0 \rightarrow$$

$$x_g = \Sigma m_i x_i / \Sigma m_i = \Sigma m_i x_i / M$$

- Calculation applies to all three body axes: x, y, z
- x can be referenced to any point, e.g., bow, waterline, geometric center, etc.
- “Enclosed” water has to be included in the mass if we are talking about inertia

Center of Buoyancy

A similar differential approach with *displaced mass*:

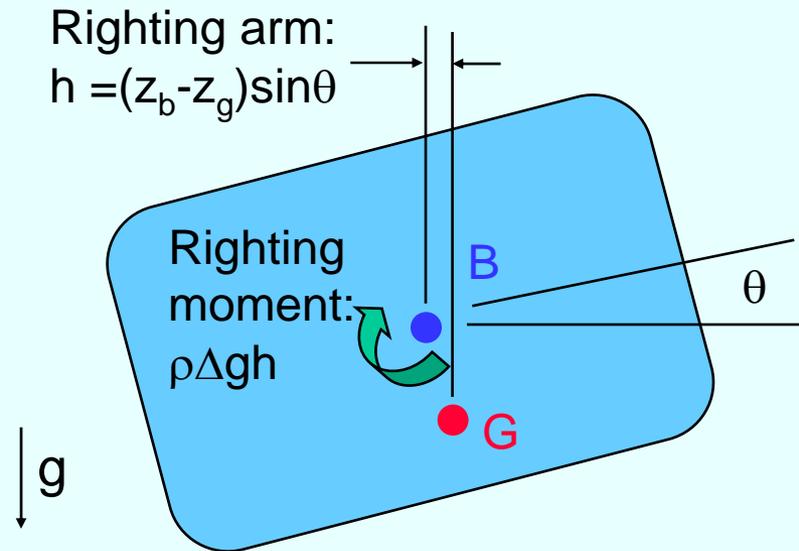
$$\mathbf{x}_b = \Sigma \Delta_i \mathbf{x}_i / \Delta, \quad \text{where } \Delta_i \text{ is incremental volume,} \\ \Delta \text{ is total volume}$$

Center of buoyancy is the same as the center of displaced volume: it doesn't matter what is inside the outer skin, or how it is arranged.



Calculating trim of a flooded vehicle: Use in-water weights of the components, including the water (whose weight is then zero and can be ignored). The calculation gives the center of in-water weight.

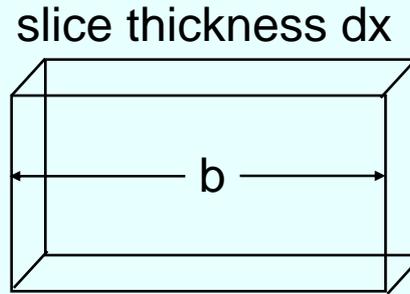
- For a submerged body, a sufficient condition for stability is that z_b is above z_g .



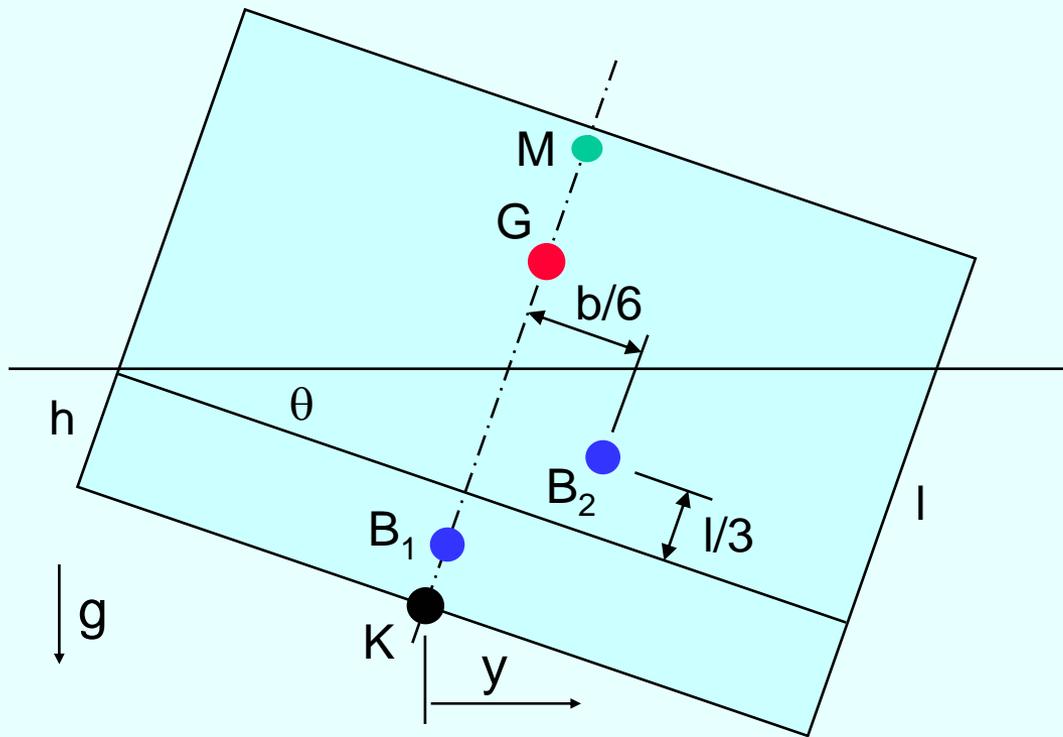
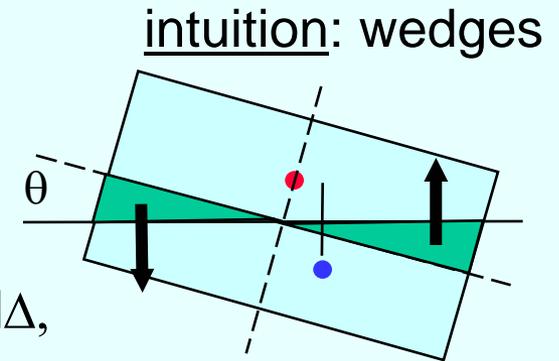
Make $(z_b - z_g)$ large \rightarrow the “spring” is large and:

- Response to an initial heel angle is fast (uncomfortable?)
- Wave or loading disturbances don't cause unacceptably large motions
- But this is also a spring-mass system, that will oscillate unless adequate damping is used, e.g., sails, anti-roll planes, etc.

- In most surface vessels, righting stability is provided by the *waterplane area*.



submerged volume $d\Delta$,
 $d\Delta = A dx$



RECTANGULAR SECTION

Geometry:

$$d\Delta/dx = bh + bl/2 \quad \text{or}$$

$$h = (d\Delta/dx - bl/2) / b$$

$$l = b \tan\theta$$

Vertical forces:

$$dF_G = -\rho g d\Delta \quad (\text{no shear})$$

$$dF_{B_1} = \rho g b h dx$$

$$dF_{B_2} = \rho g b l dx / 2$$

Moment arms:

$$y_G = KG \sin \theta \quad ; \quad y_{B1} = h \sin \theta / 2 \quad ; \quad y_{B2} = (h + l/3) \sin \theta + b \cos \theta / 6$$

Put all this together into a net moment (positive anti-clockwise):

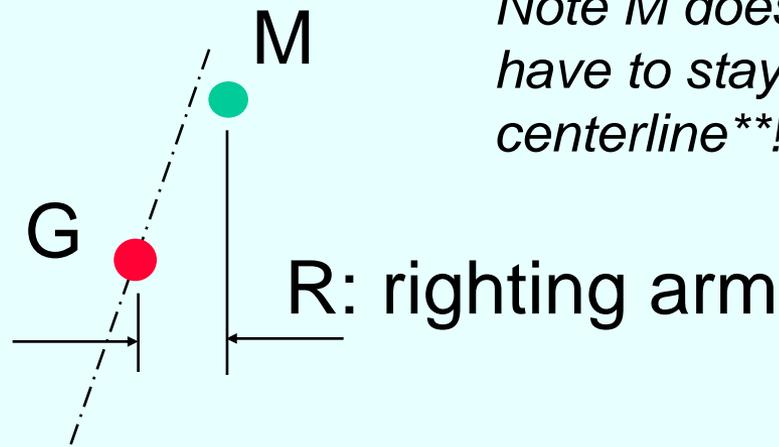
$$dM/\rho g = -KG \, d\Delta \sin\theta + bh^2 \, dx \sin\theta / 2 + \quad \text{(valid until the corner comes out of the water)}$$
$$b \, l \, dx \left[(h+l/3) \sin\theta + b \cos\theta / 6 \right] / 2$$

Linearize ($\sin\theta \sim \tan\theta \sim \theta$), and keep only first-order terms (θ):

$$dM / \rho g \, d\Delta = [-KG + h / 2 + b^2 / 12 h] \theta$$
$$= [-KG + A / 2 b + b^3 / 12 A] \theta$$

For this rectangular slice, *the sum $[h / 2 + b^2 / 12 h]$ must exceed the distance KG for stability.* This sum is called KM – the distance from the keel up to the “virtual” buoyancy center M . M is the METACENTER, and it is as if the block is hanging from M !

$-KG + KM = GM$: the METACENTRIC HEIGHT



Note M does not have to stay on the centerline**!

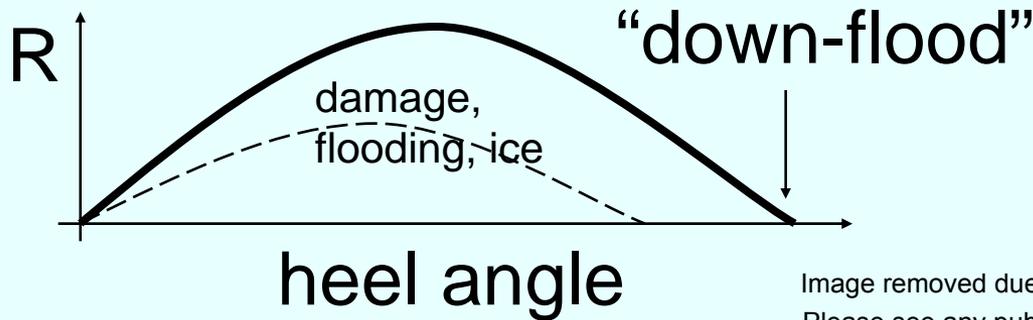


Image removed due to copyright restrictions.

Please see any publicity images for *The Perfect Storm*, such as <http://www.imdb.com/media/rm1041734656/tt0177971>

How much GM is enough?
Around 2-3m in a big boat

The Perfect Storm

Considering the Entire Vessel...

Transverse (or roll) stability is calculated using the same moment calculation extended on the length:

Total Moment = Integral on Length of $dM(x)$, where (for a vessel with all rectangular cross-sections)

$$dM(x) = \rho g [-KG(x) A(x) + A^2(x) / 2 b(x) + b^3(x) / 12] dx \theta$$

First term: Same as $-\rho g KG \Delta$, if Δ is the ship's submerged volume, and KG is the value referencing the whole vessel

Second term: Significant if $d > b$ (equivalent to $h^2 b / 2$)

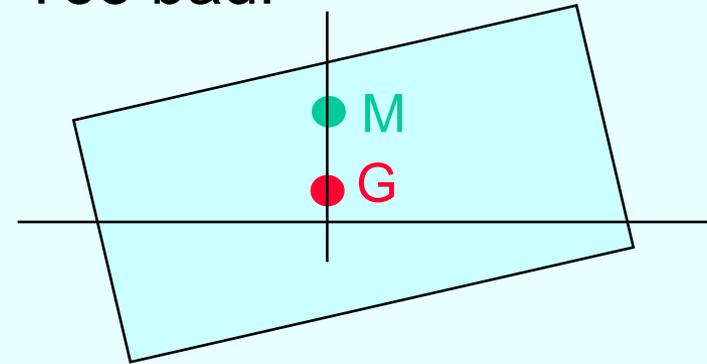
Third term: depends only on beam – dominant for most monohulls

Longitudinal (or pitch) stability is similarly calculated, but it is usually secondary, since the waterplane area is very long \rightarrow very high GM

Weight Distribution and Trim

- At zero speed, and with no other forces or moments, the vessel has B (submerged) or M (surface) *directly above* G.

Too bad!



For port-stbd symmetric hulls, keep G on the centerline using a tabulation of component masses and their centroid locations in the hull, i.e., $\sum m_i y_i = 0$

Longitudinal trim should be zero relative to center of waterplane area, in the loaded condition.

Pitch trim may be affected by forward motion, but difference is usually only a few degrees.

Rotational Dynamics Using the Centroid

Equivalent to $F = ma$ in linear case is

$$T = J_o * d^2\theta / dt^2$$

where T is the sum of acting torques in roll

J_o is the rotary moment of inertia in roll,
referenced to some location O

θ is roll angle (radians)

J written in terms of incremental masses m_i :

$$J_o = \sum m_i (y_i - y_o)^2 \quad \text{OR} \quad J_g = \sum m_i (y_i - y_g)^2$$

J written in terms of component masses m_i and their own moments of inertia J_i (by the parallel axis theorem) :

$$J_g = \sum m_i (y_i - y_g)^2 + \sum J_i$$

The y_i 's give position of the centroid of each body, and J_i 's are referenced to those centroids

What are the acting torques T ?

- Buoyancy righting moment – metacentric height
- Dynamic loads on the vessel – e.g., waves, wind, movement of components, sloshing
- Damping due to keel, roll dampers, etc.
- Torques due to roll control actuators

An instructive case of damping D , metacentric height GM :

$$J d^2\theta / dt^2 = - D d\theta / dt - GM \rho g \Delta \theta \quad \text{OR}$$

$$J d^2\theta / dt^2 + D d\theta / dt + GM \rho g \Delta \theta = 0$$

$$d^2\theta / dt^2 + a d\theta / dt + b\theta = 0$$

$$d^2\theta / dt^2 + 2 \zeta \omega_n d\theta / dt + \omega_n^2 \theta = 0$$

A second-order stable system \rightarrow Overdamped or oscillatory response from initial conditions

Homogeneous Underdamped Second-Order Systems

$$x'' + ax' + bx = 0; \quad \text{write as} \quad x'' + 2\zeta\omega_n x' + \omega_n^2 x = 0$$

Let $x = X e^{st} \rightarrow$

$$(s^2 + 2\zeta\omega_n s + \omega_n^2) e^{st} = 0 \quad \text{OR} \quad s^2 + 2\zeta\omega_n s + \omega_n^2 = 0 \rightarrow$$
$$s = \frac{-2\zeta\omega_n \pm \sqrt{4\zeta^2\omega_n^2 - 4\omega_n^2}}{2}$$
$$= \omega_n[-\zeta \pm \sqrt{\zeta^2 - 1}] \quad \text{from quadratic equation}$$

s_1 and s_2 are complex conjugates if $\zeta < 1$, in this case:

$$s_1 = -\omega_n\zeta + i\omega_d, \quad s_2 = -\omega_n\zeta - i\omega_d \quad \text{where} \quad \omega_d = \omega_n \sqrt{1 - \zeta^2}$$

Recalling $e^{r+i\theta} = e^r (\cos\theta + i \sin\theta)$, we have

$$x = e^{-\zeta\omega_n t} \left[(X_1^r + iX_1^i)(\cos\omega_d t + i\sin\omega_d t) + (X_2^r + iX_2^i)(\cos\omega_d t - i\sin\omega_d t) \right] \quad \text{AND}$$

$$\mathbf{x}' = -\zeta\omega_n\mathbf{x} + \omega_d e^{-\zeta\omega_n t} \left[(\mathbf{X}_1^r + i\mathbf{X}_1^i)(-\sin\omega_d t + i\cos\omega_d t) + (\mathbf{X}_2^r + i\mathbf{X}_2^i)(-\sin\omega_d t - i\cos\omega_d t) \right]$$

Consider initial conditions $\mathbf{x}'(0) = 0$, $\mathbf{x}(0) = 1$:

$$\mathbf{x}(t=0) = 1 \text{ means } \begin{array}{l} \mathbf{X}_1^r + \mathbf{X}_2^r = 1 \quad (\text{real part}) \text{ and} \\ \mathbf{X}_1^i + \mathbf{X}_2^i = 0 \quad (\text{imaginary part}) \end{array}$$

$$\mathbf{x}'(t=0) = 0 \text{ means } \begin{array}{l} \mathbf{X}_1^r - \mathbf{X}_2^r = 0 \quad (\text{imaginary part}) \text{ and} \\ -\zeta\omega_n + \omega_d(\mathbf{X}_2^i - \mathbf{X}_1^i) = 0 \quad (\text{real part}) \end{array}$$

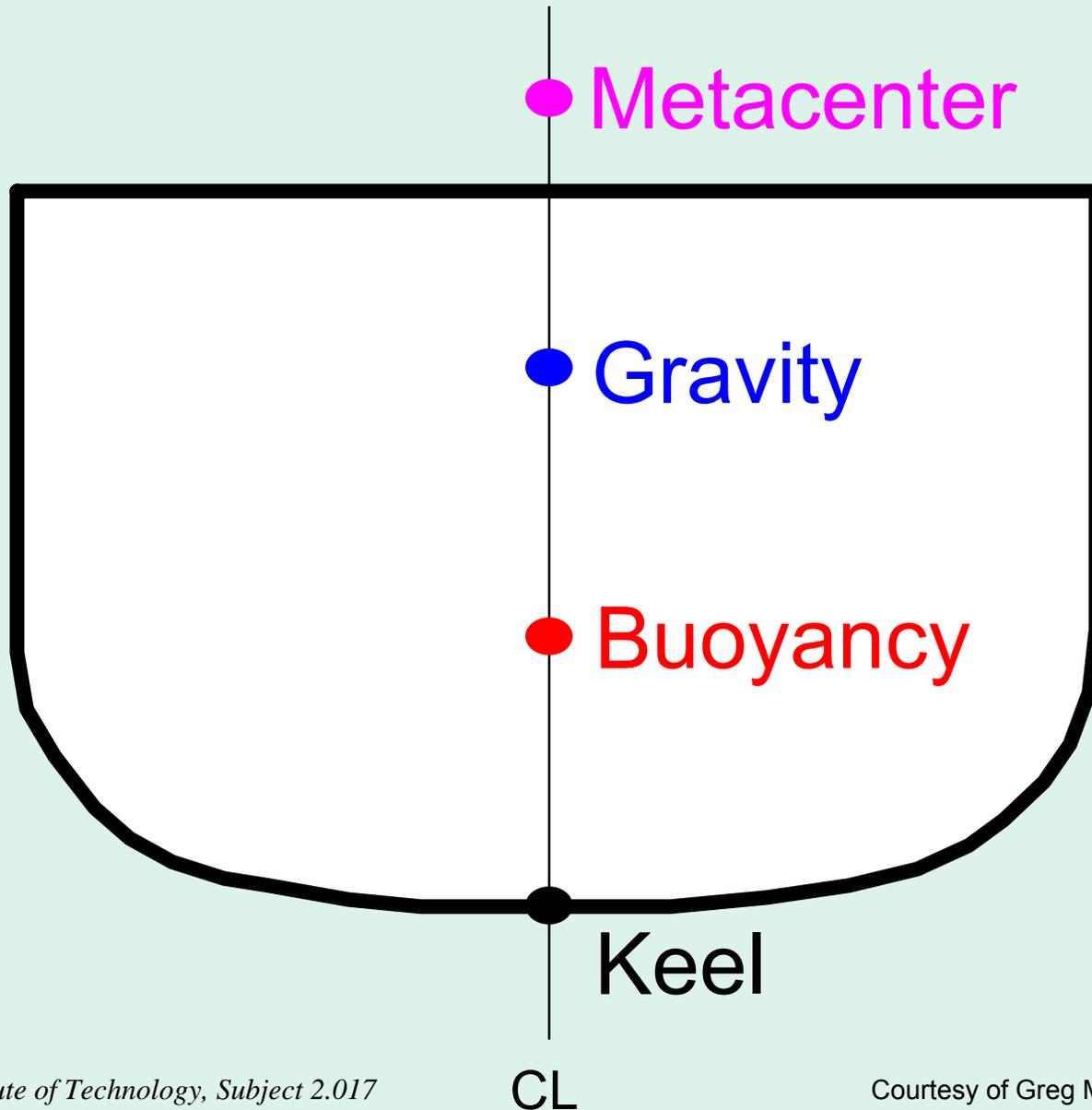
Combine these and we find that

$$\begin{array}{l} \mathbf{X}_1^r = \mathbf{X}_2^r = 1/2 \\ \mathbf{X}_1^i = -\mathbf{X}_2^i = -\zeta\omega_n / 2\omega_d \end{array}$$

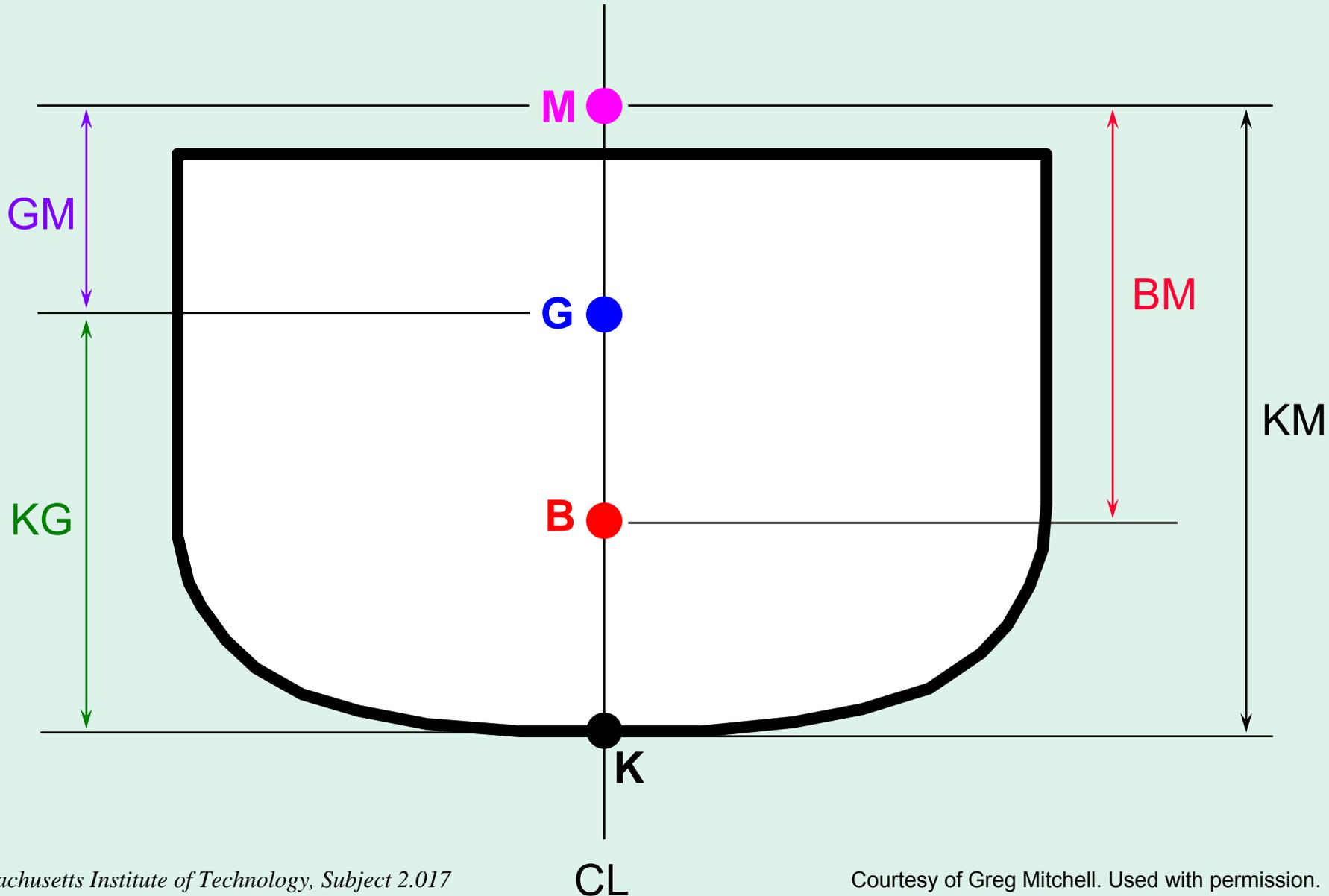
Plug into the solution for \mathbf{x} and do some trig:

$$\mathbf{x} = e^{-\zeta\omega_n t} \sin(\omega_d t + \mathbf{k}) / \text{sqrt}(1-\zeta^2), \text{ where } \mathbf{k} = \text{atan}(\omega_d/\zeta\omega_n)$$

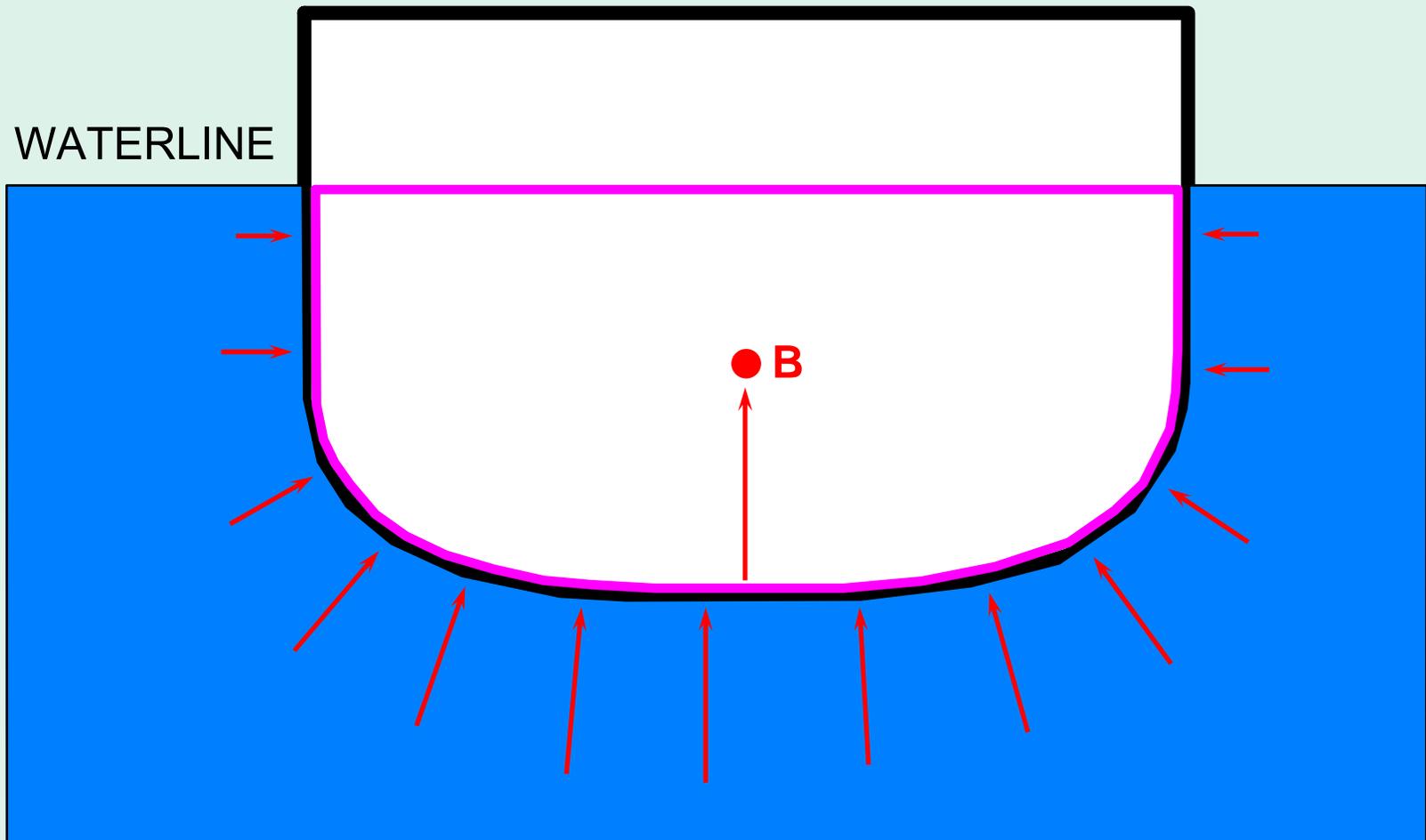
STABILITY REFERENCE POINTS



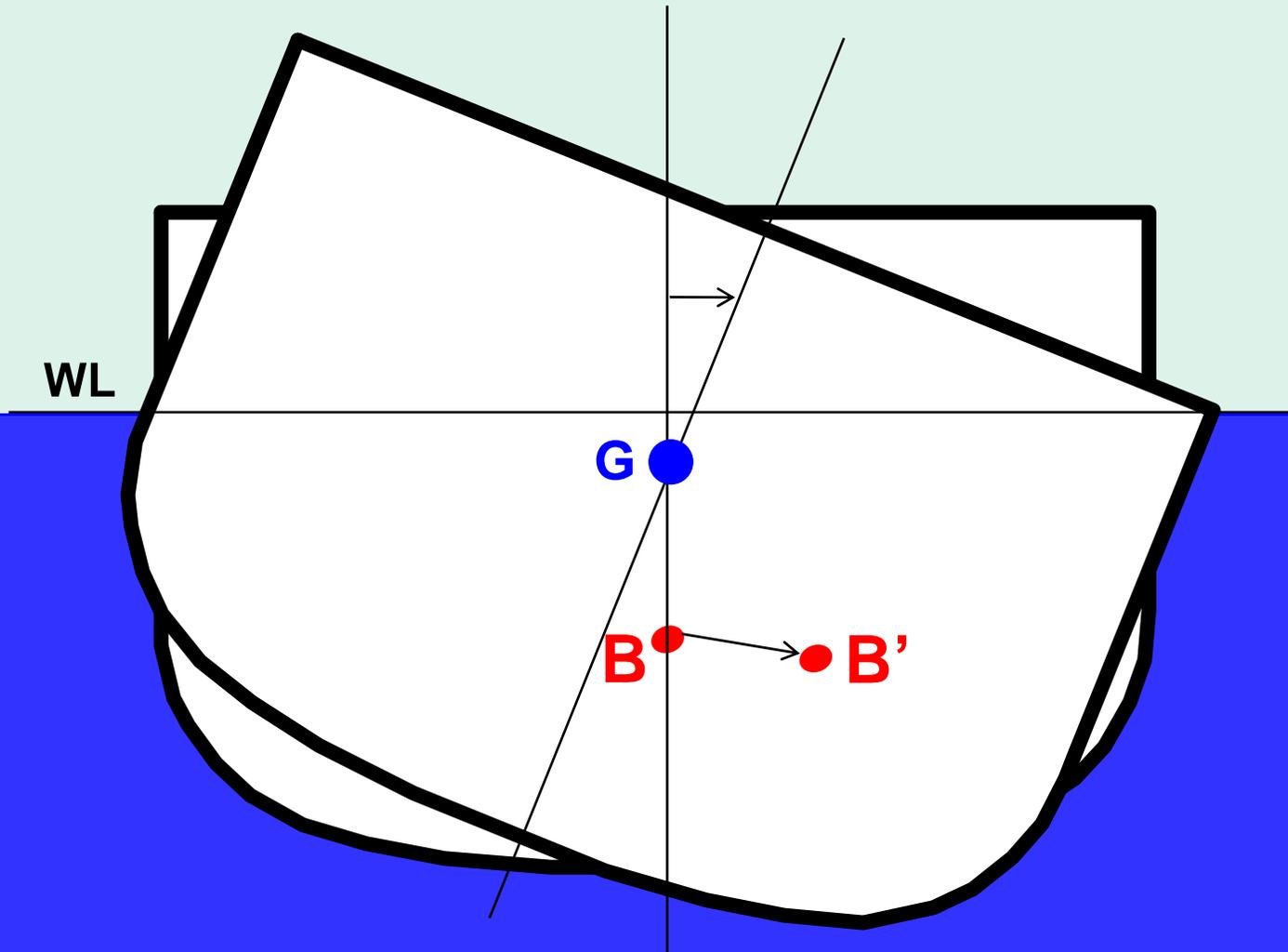
LINEAR MEASUREMENTS IN STABILITY



THE CENTER OF BUOYANCY

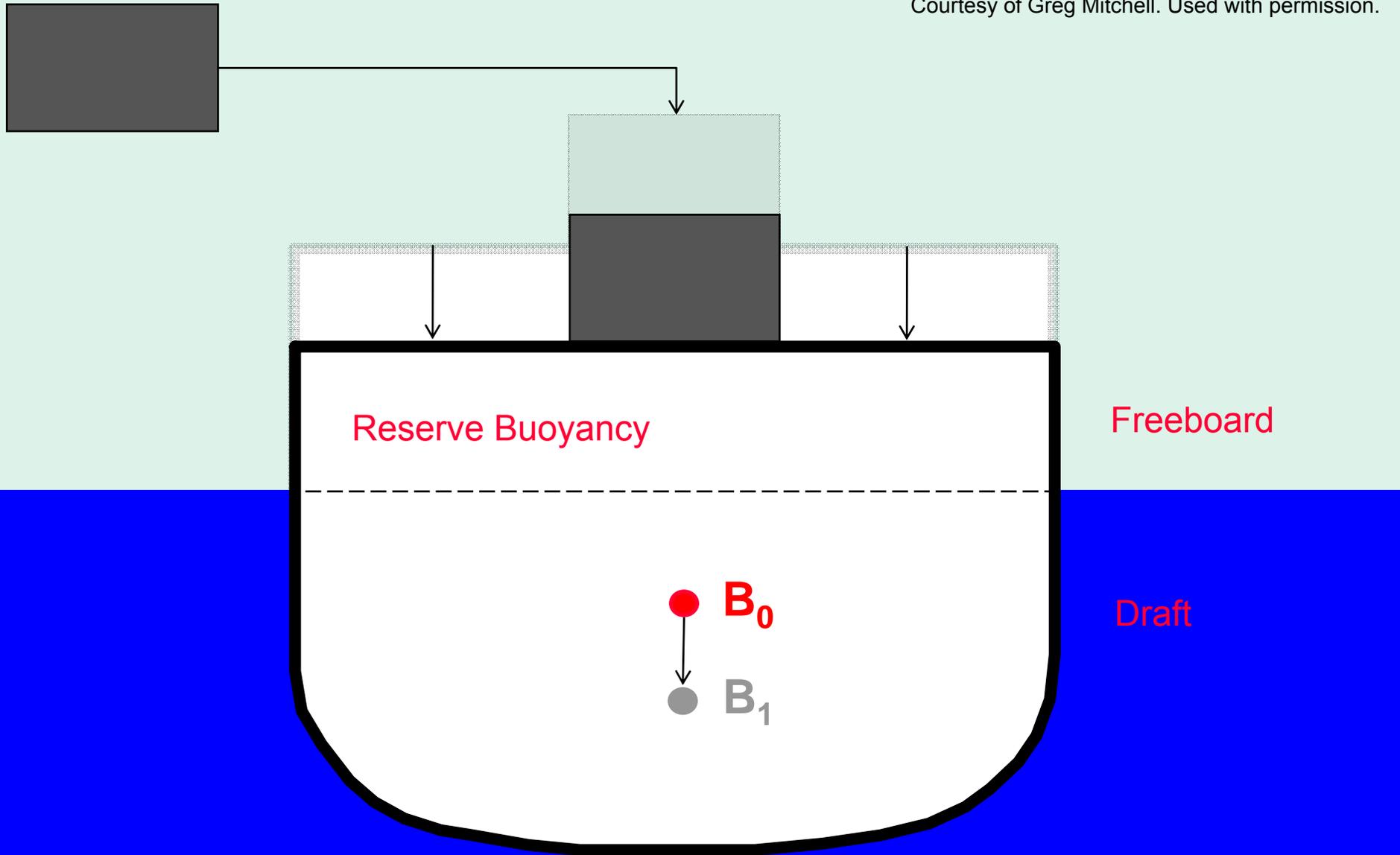


CENTER OF BUOYANCY



CENTER OF BUOYANCY

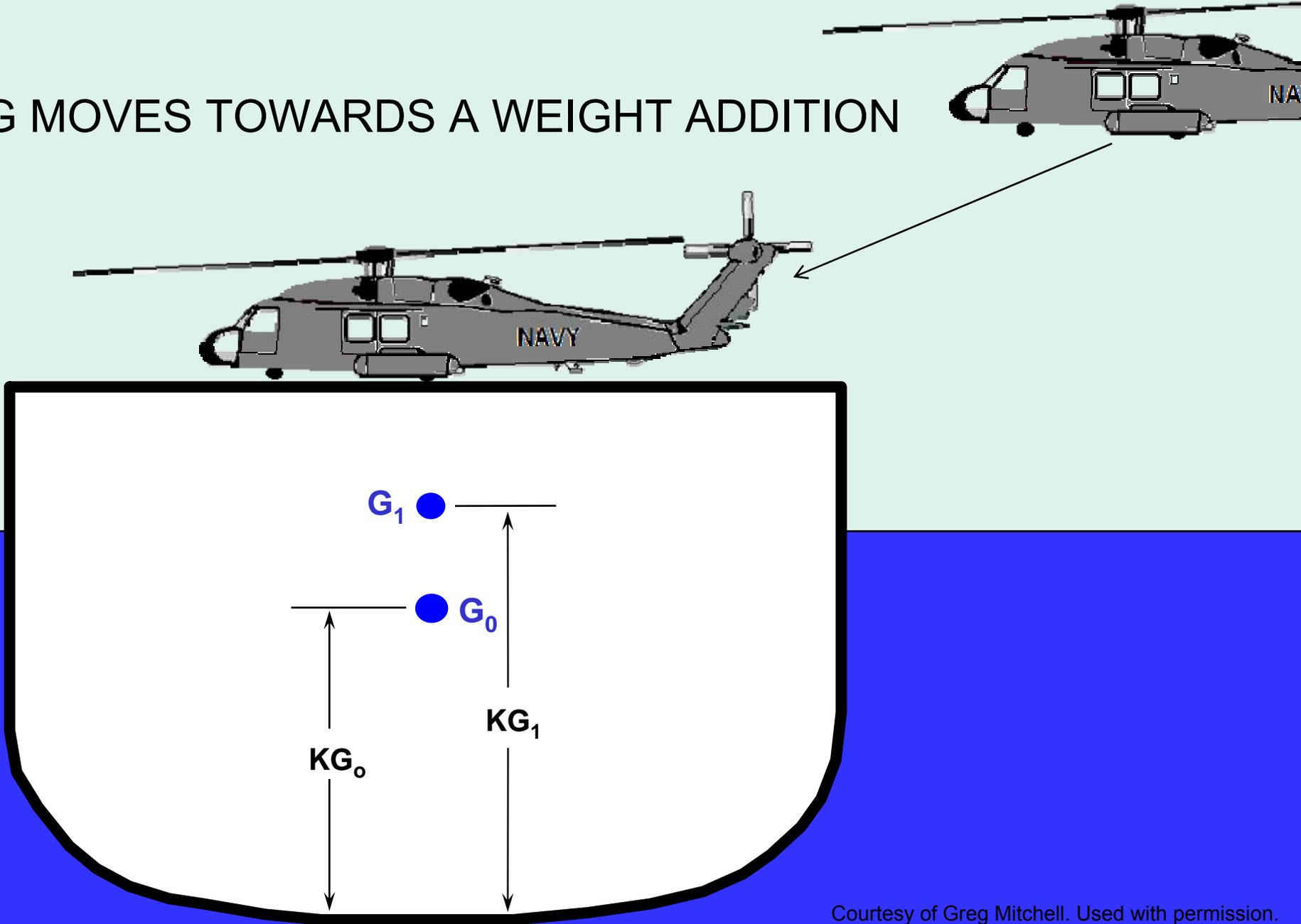
Courtesy of Greg Mitchell. Used with permission.



- The freeboard and reserve buoyancy will also change

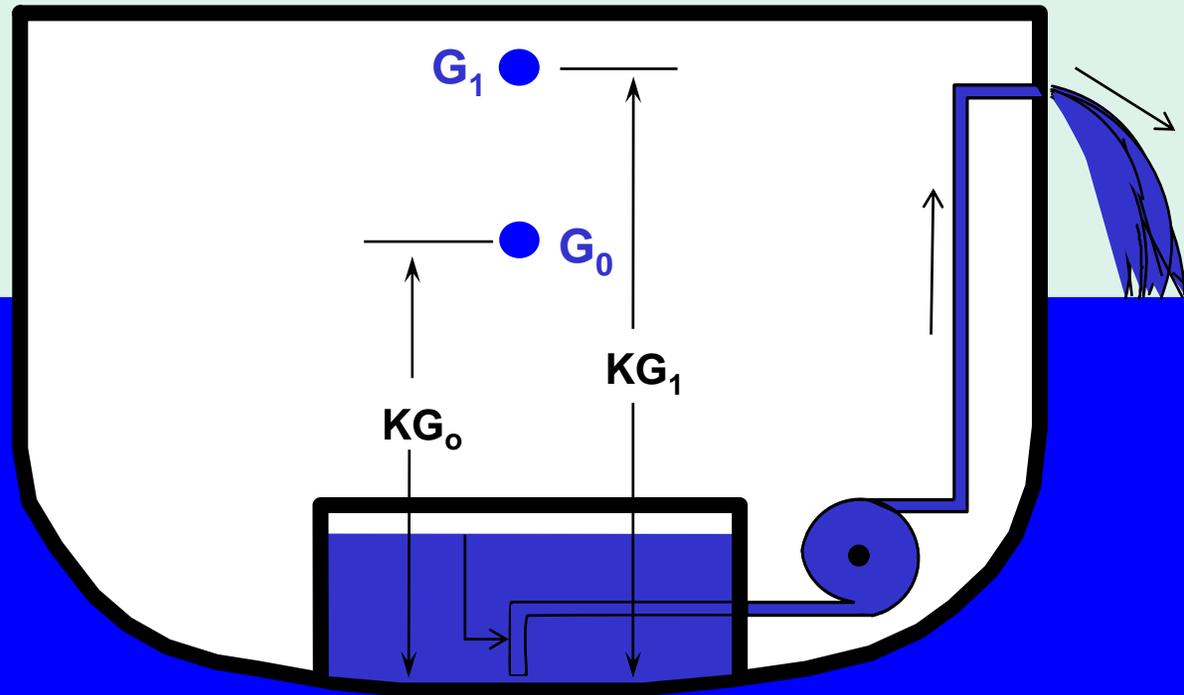
MOVEMENTS IN THE CENTER OF GRAVITY

G MOVES TOWARDS A WEIGHT ADDITION

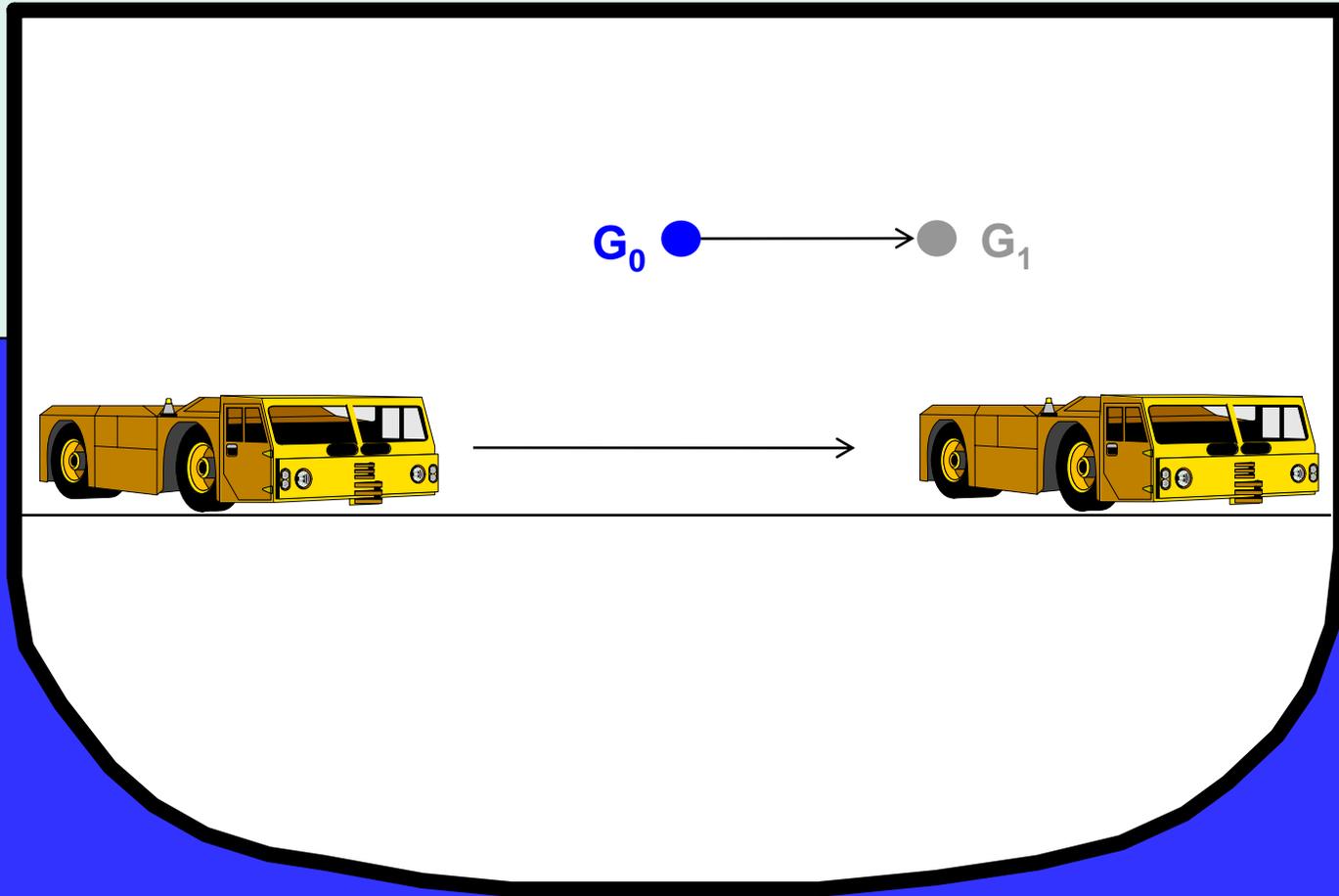


MOVEMENTS IN THE CENTER OF GRAVITY

G MOVES AWAY FROM A WEIGHT REMOVAL

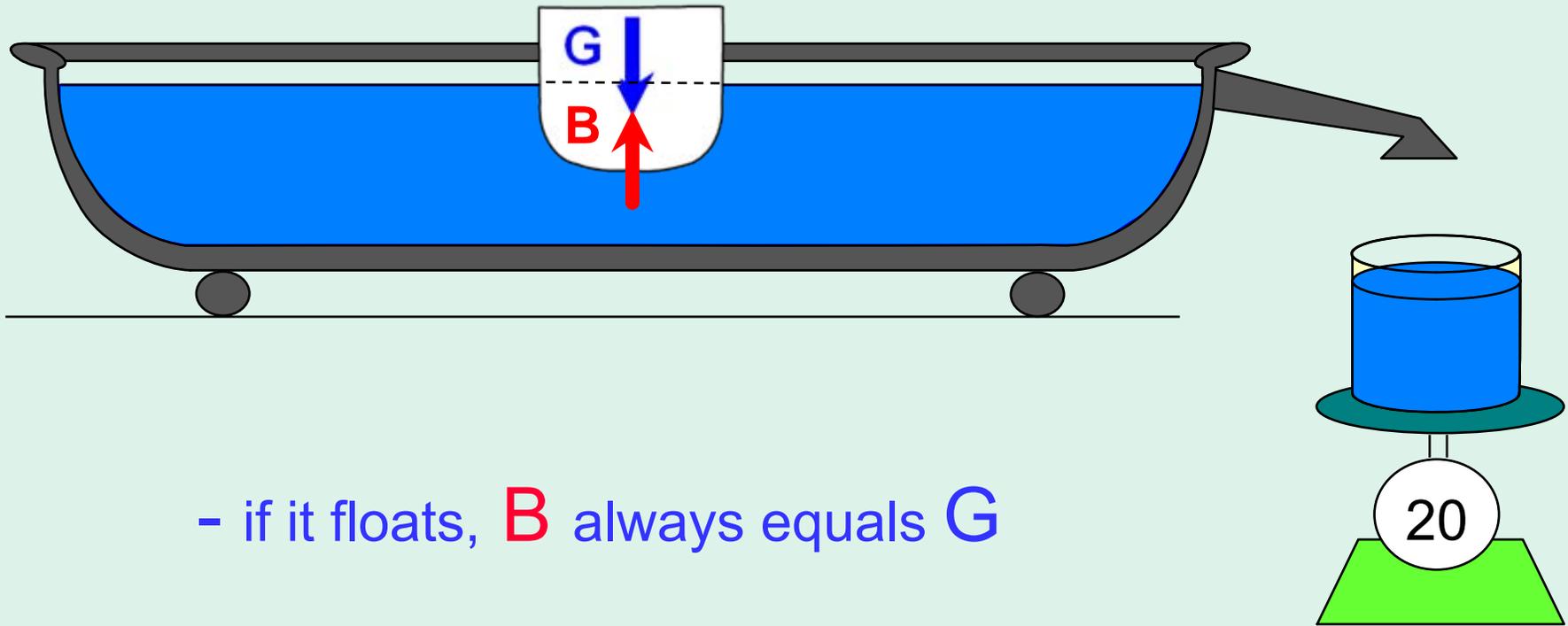


MOVEMENTS IN THE CENTER OF GRAVITY



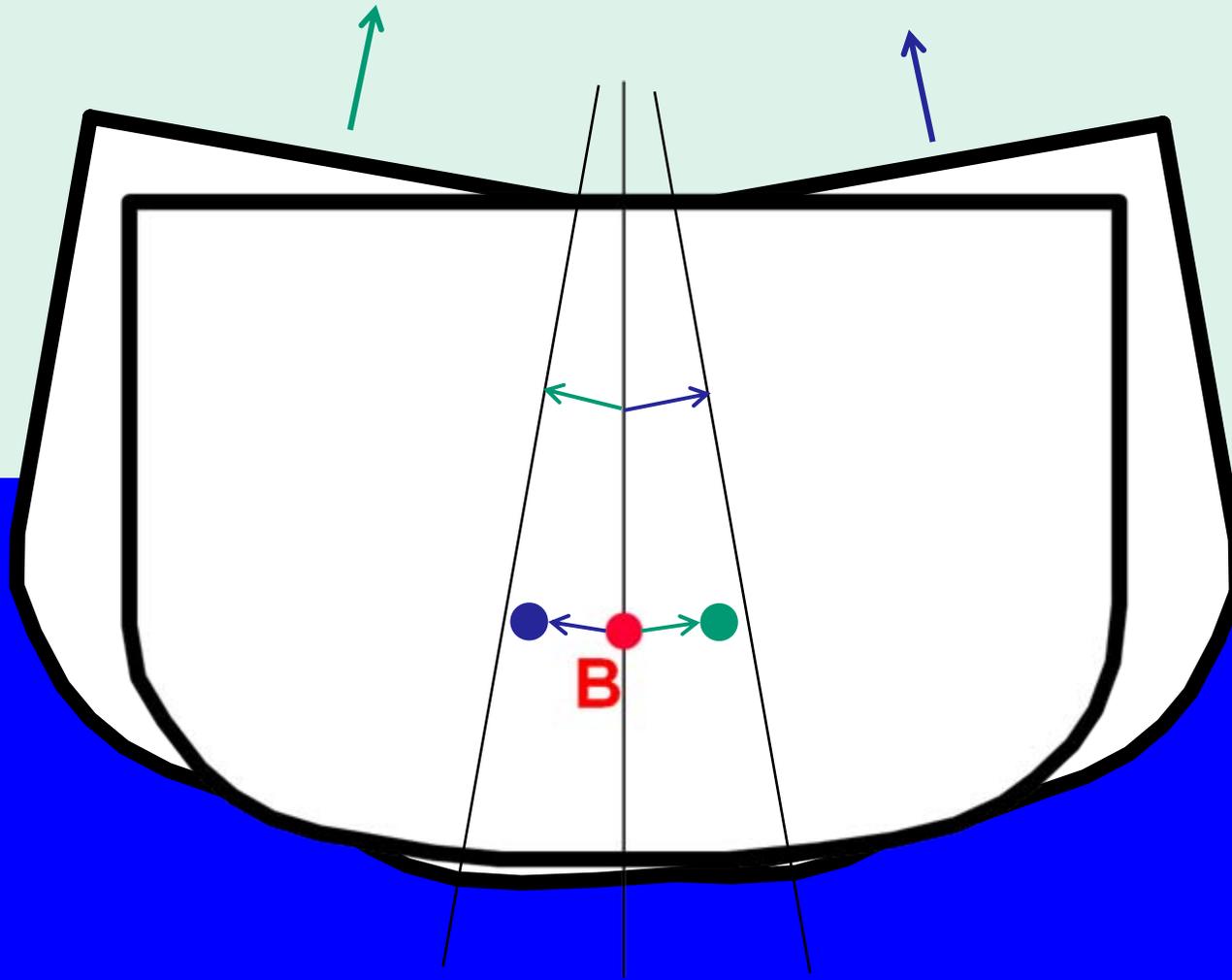
G MOVES IN THE DIRECTION OF A WEIGHT SHIFT

DISPLACEMENT = SHIP'S WEIGHT

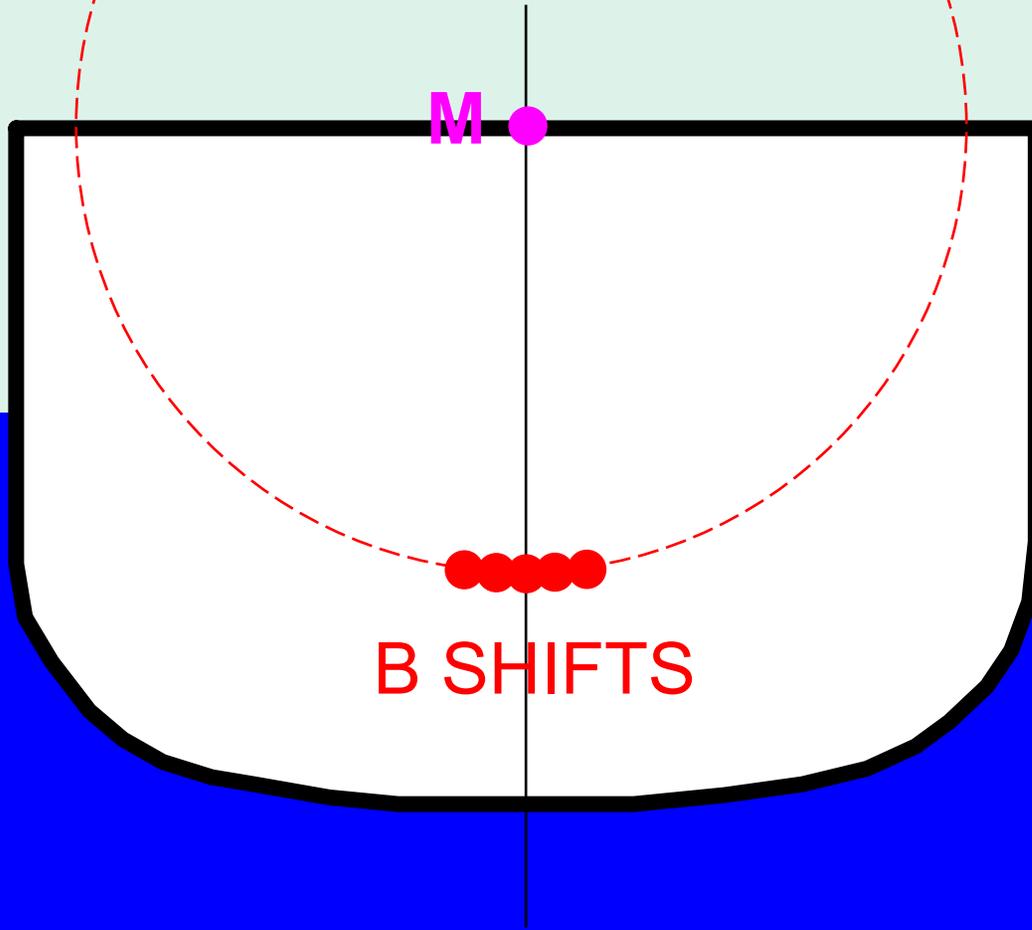


- if it floats, **B** always equals **G**

METACENTER

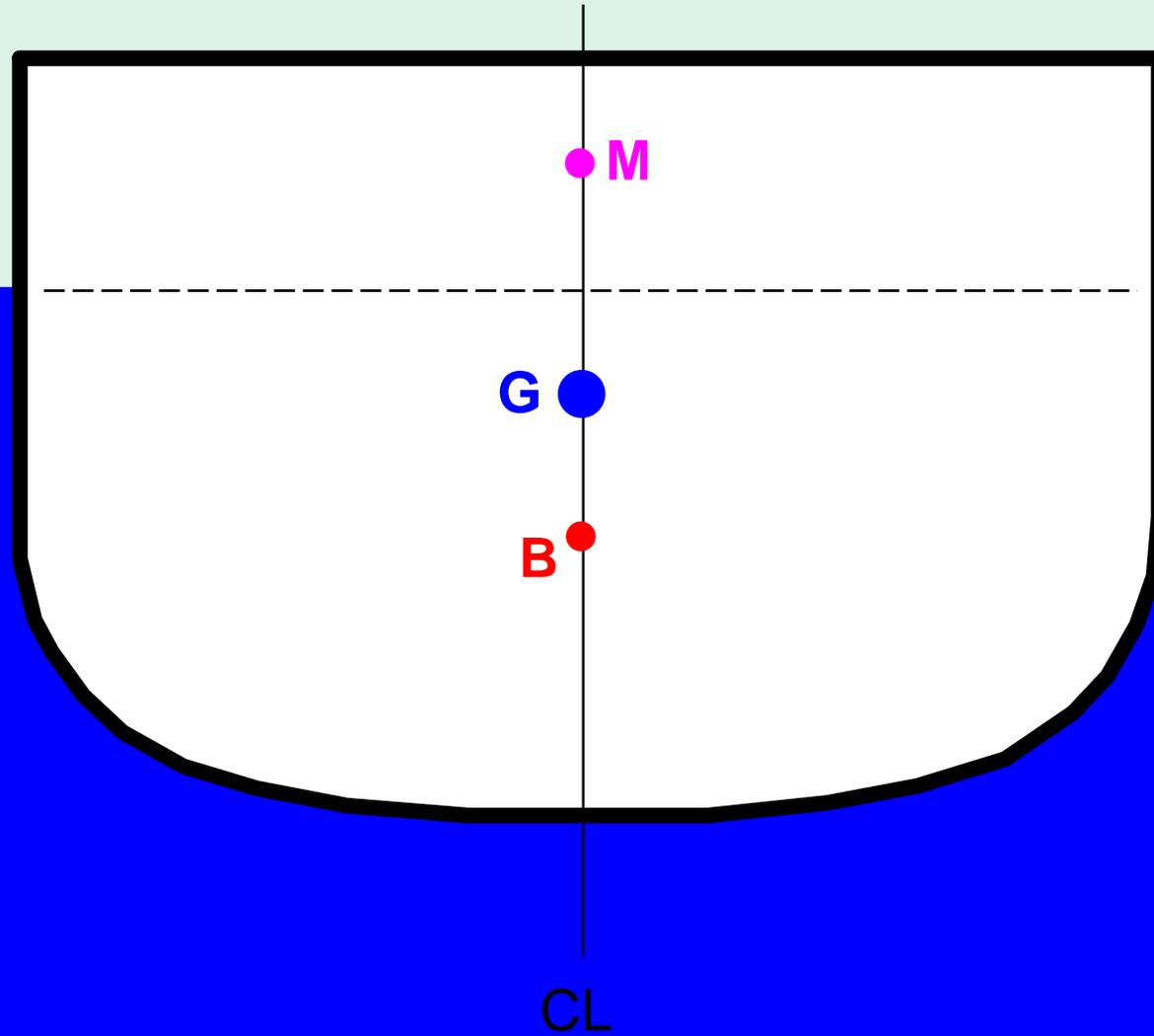


METACENTER

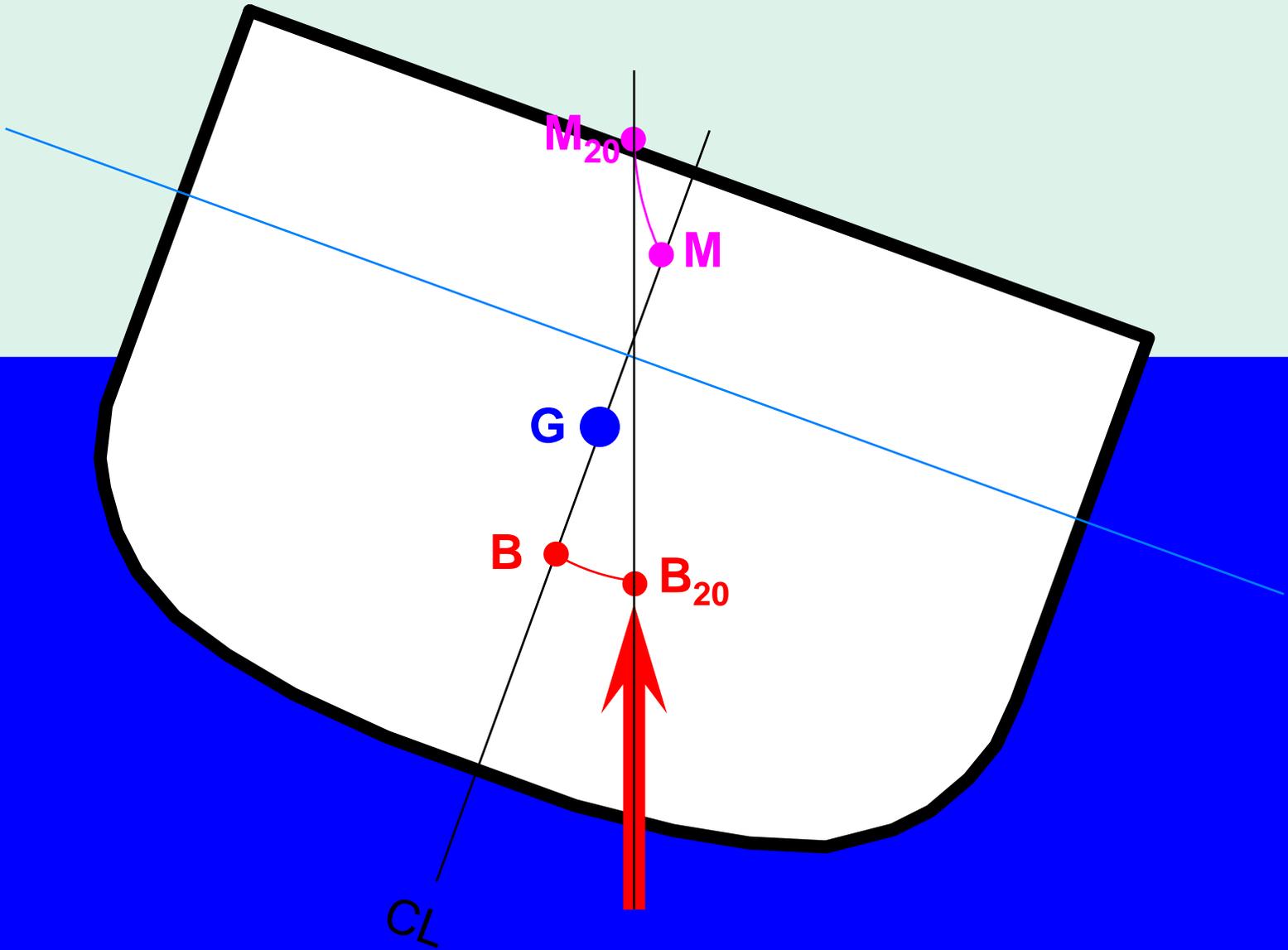


+GM

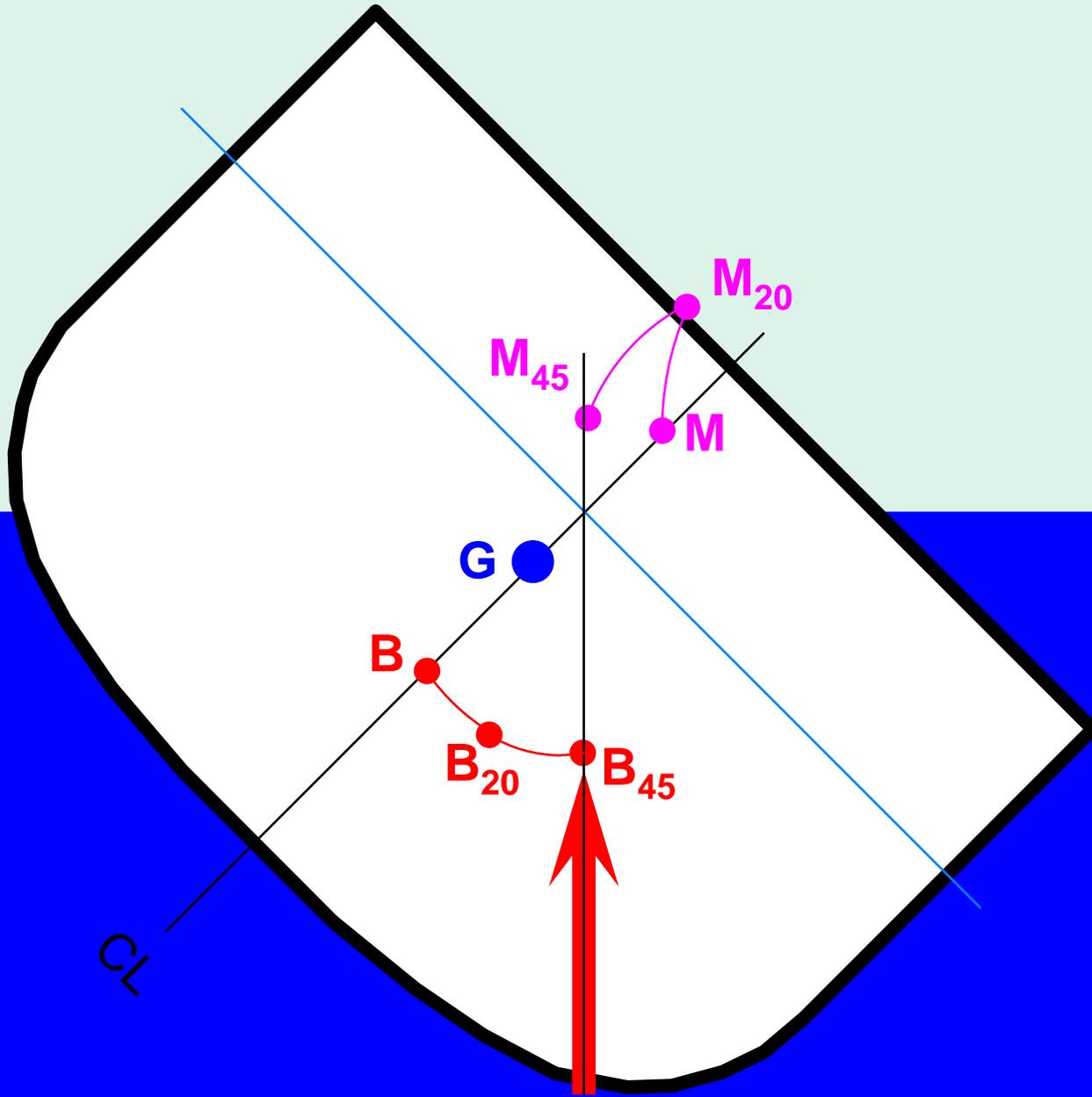
$0^\circ - 7/10^\circ$



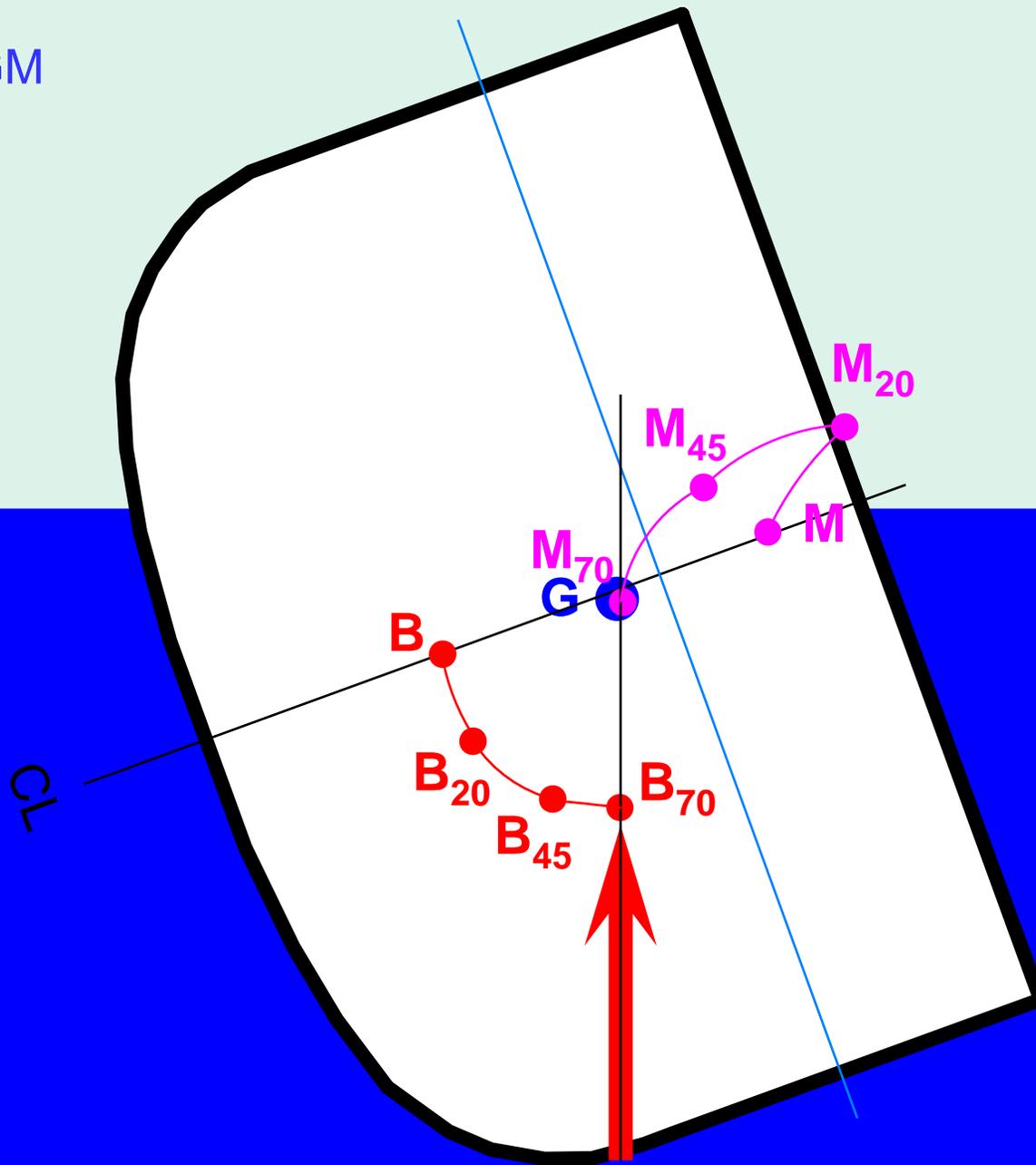
+GM



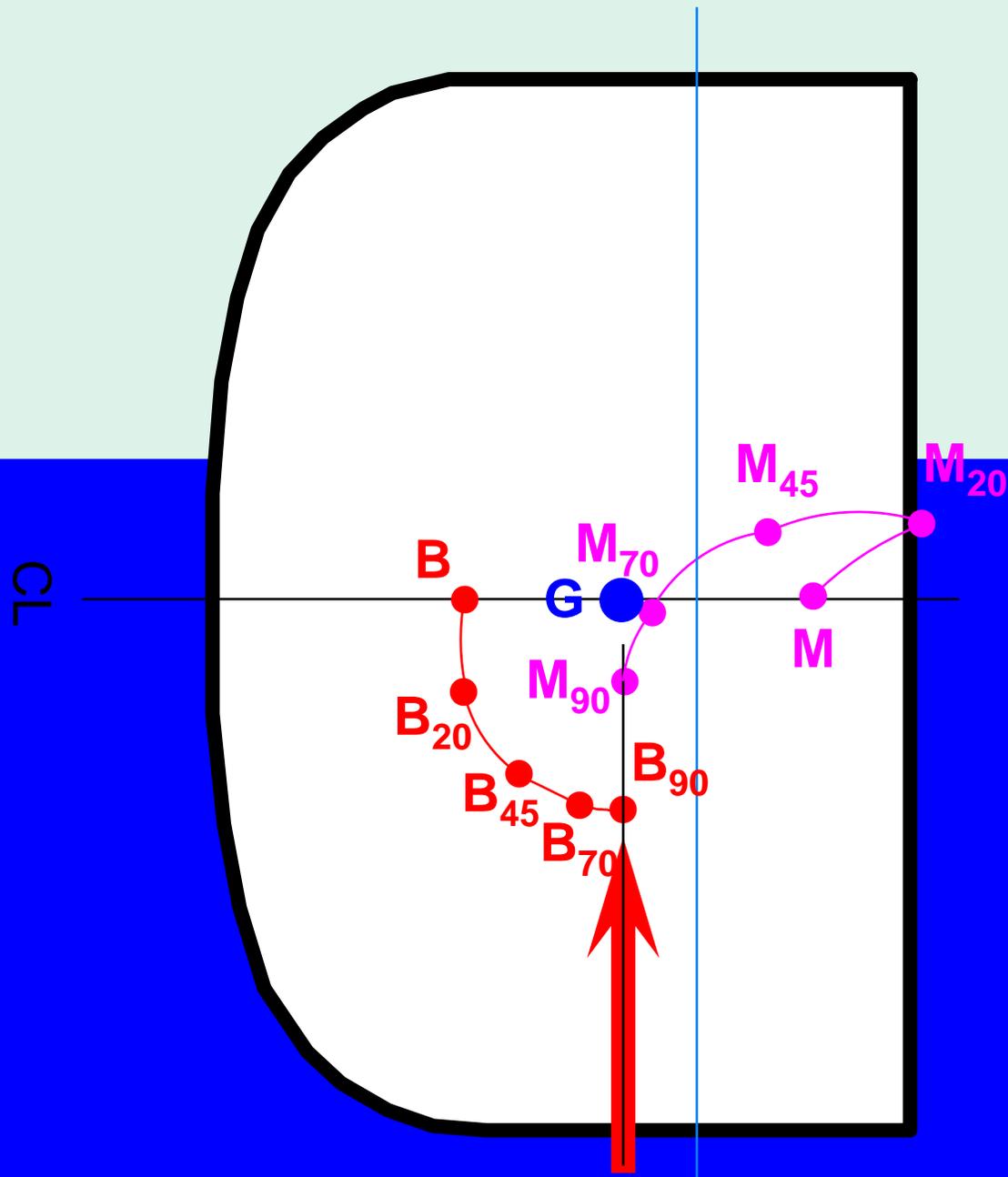
+GM



neutral GM



-GM

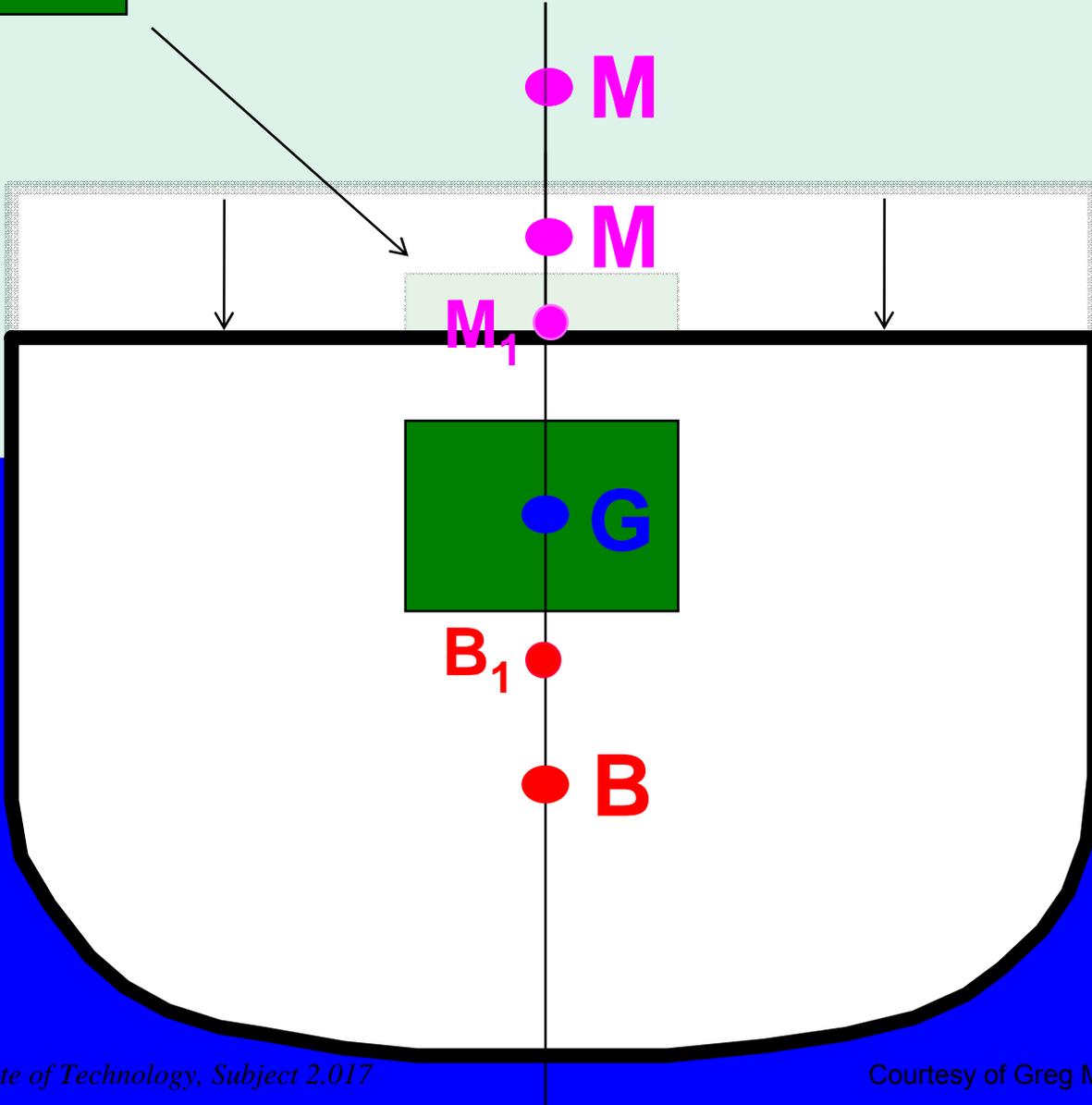
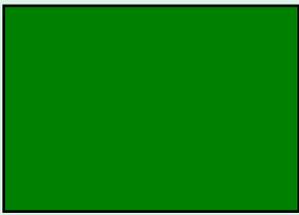


MOVEMENTS OF THE METACENTER

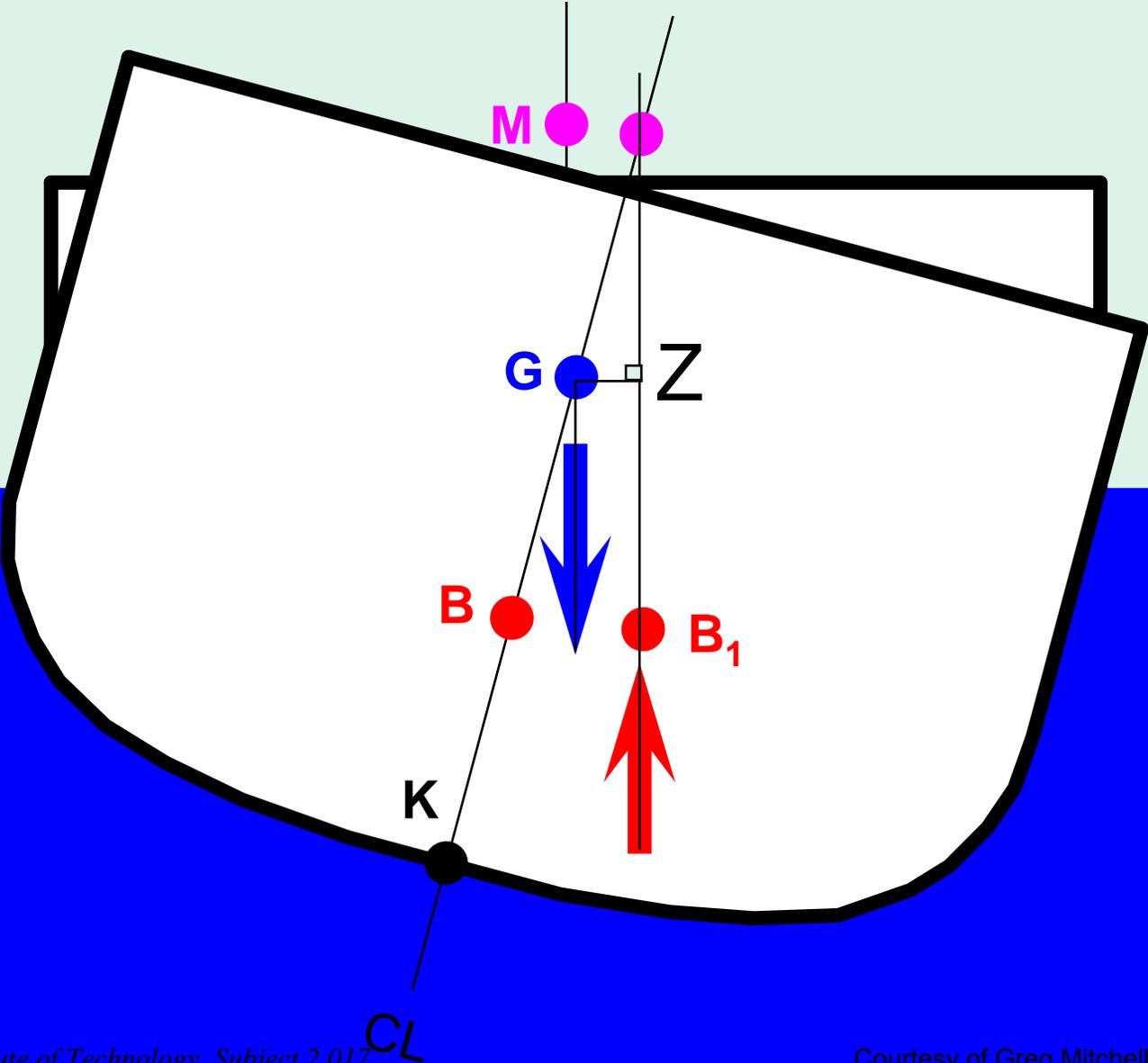
THE METACENTER WILL CHANGE POSITIONS IN THE VERTICAL PLANE WHEN THE SHIP'S DISPLACEMENT CHANGES

THE METACENTER MOVES IAW THESE TWO RULES:

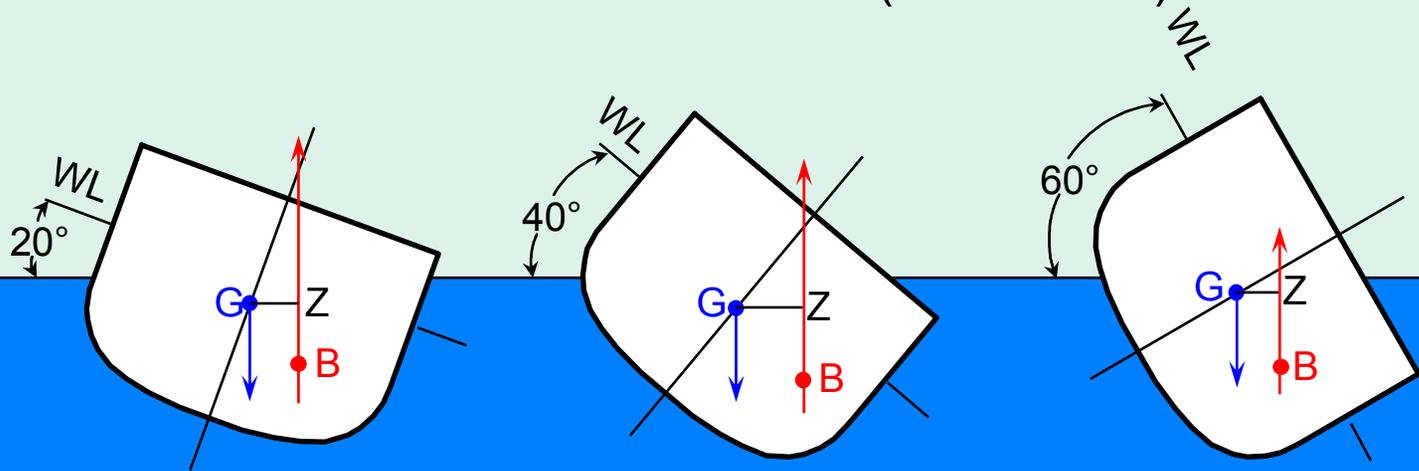
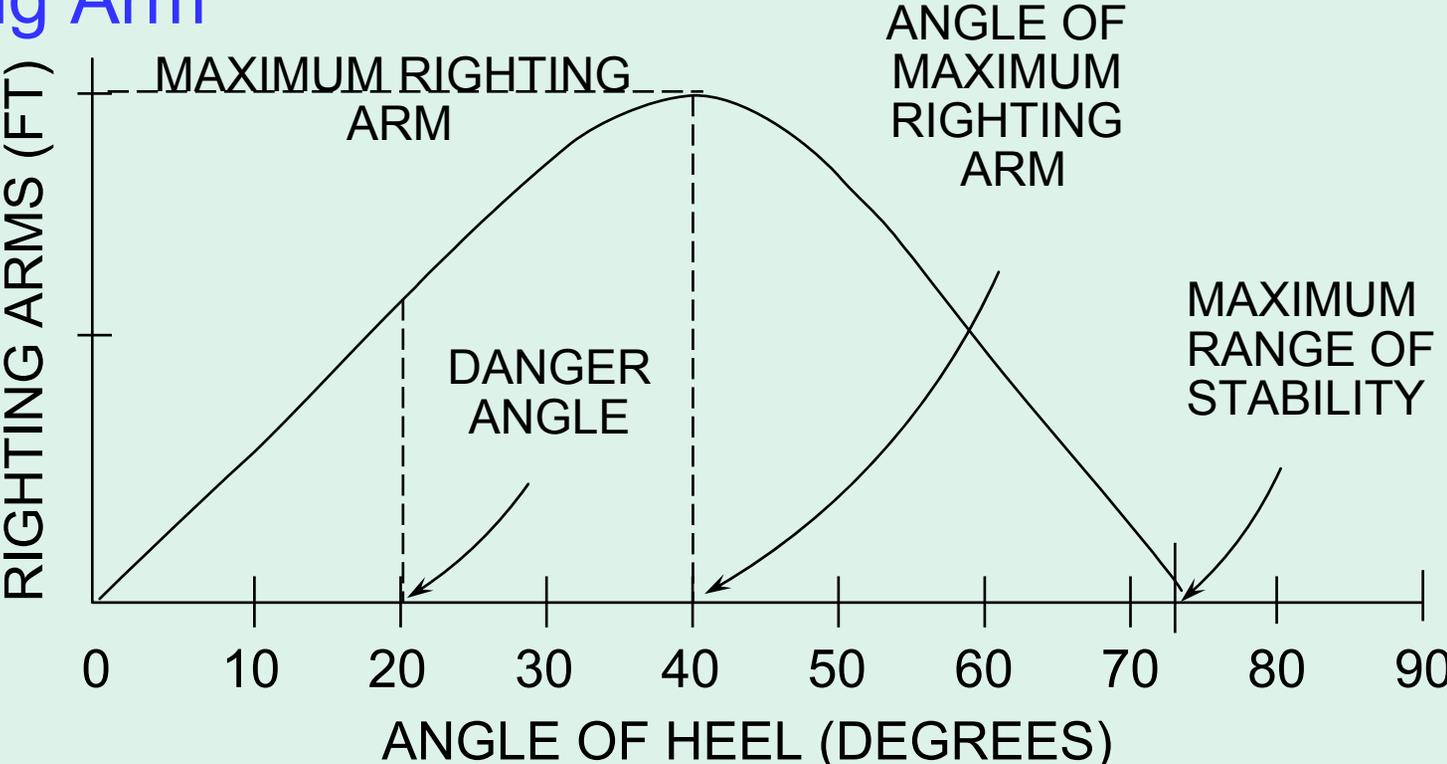
1. WHEN B MOVES UP, M MOVES DOWN.
2. WHEN B MOVES DOWN, M MOVES UP.



Righting Arm



Righting Arm

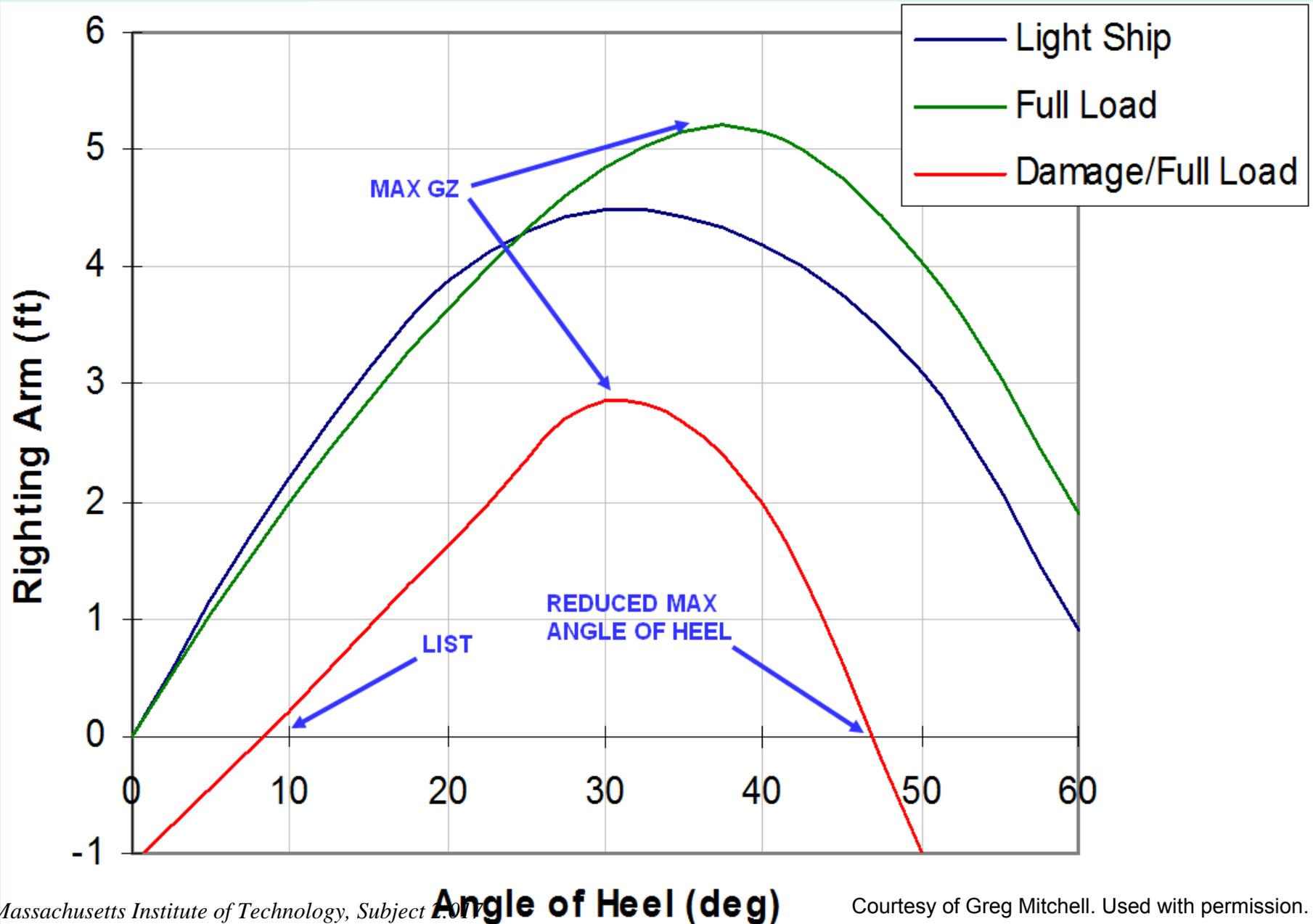


GZ = 1.4 FT

GZ = 2.0 FT

GZ = 1 FT

Righting Arm for Various Conditions



THINGS TO CONSIDER

- Effects of:
 - Weight addition/subtraction and movement
 - Ballasting and loading/unloading operations
 - Wind, Icing
 - Damage stability
 - result in an adverse movement of G or B
 - sea-keeping characteristics will change
 - compensating for flooding (ballast/completely flood a compartment)
 - maneuvering for seas/wind

References

- NSTM 079 v. I Buoyancy & Stability
- NWP 3-20.31 Ship Survivability
- Ship's Damage Control Book
- Principles of Naval Architecture v. I

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