

6.730 Physics for Solid State Applications

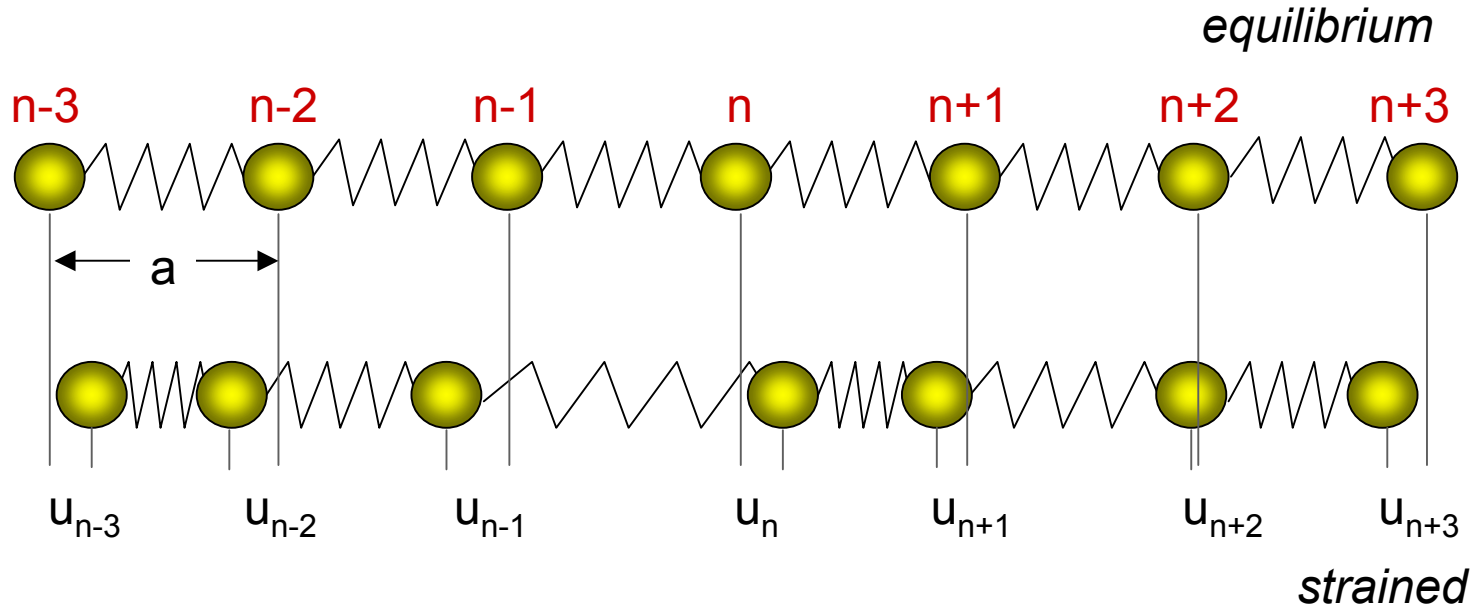
Lecture 9: Lattice Waves in 1D with Diatomic Basis

Outline

- Review Lecture 8
- 1-D Lattice with Basis
- Example of Nearest Neighbor Coupling
- Optical and Acoustic Phonon Branches

Strain in a Discrete 1-D Monatomic Lattice

General Expansion



$$F[n, t] = - \left(\frac{\partial V}{\partial u[n, t]} \right)_{eq} = 0$$

$$V(\{u[i, t]\}) = V_0 + \sum_{m=-\infty}^{\infty} \left(\frac{\partial V}{\partial u[m, t]} \right)_{eq} u[m, t]$$

$$+ \frac{1}{2} \sum_{n=-\infty}^{\infty} \sum_{m=-\infty}^{\infty} u[n, t] \left(\frac{\partial^2 V}{\partial u[n, t] \partial u[m, t]} \right)_{eq} u[m, t] + \dots$$

Equations of Motion for Lattice Atoms

$$V(\{u[i, t]\}) = V_o + \frac{1}{2} \sum_{n=-\infty}^{\infty} \sum_{m=-\infty}^{\infty} u[n, t] \left(\frac{\partial^2 V}{\partial u[n, t] \partial u[m, t]} \right)_{\text{eq}} u[m, t] + \dots$$

Harmonic Matrix: $\tilde{D}(n, m) = \left(\frac{\partial^2 V}{\partial u[n, t] \partial u[m, t]} \right)_{\text{eq}}$

$$V(\{u[i, t]\}) = V_o + \frac{1}{2} \sum_{n=-\infty}^{\infty} \sum_{m=-\infty}^{\infty} u[n, t] \tilde{D}(n, m) u[m, t]$$

Force on the j^{th} atom (away from equilibrium)...

$$M \frac{d^2}{dt^2} u[j] = - \frac{\partial}{\partial u[j]} V(\{u[i]\}) = - \sum_{m=-\infty}^{\infty} \tilde{D}(j, m) u[m]$$

Solutions of Equations of Motion

$$M \frac{d^2}{dt^2} u[n, t] = - \sum_{m=-\infty}^{\infty} \tilde{D}(n, m) u[m, t]$$

Assuming time-harmonic solutions, converts into coupled difference equations:

$$M\omega^2 \tilde{U}[n] = \sum_{m=-\infty}^{\infty} \tilde{D}(n, m) \tilde{U}[m]$$

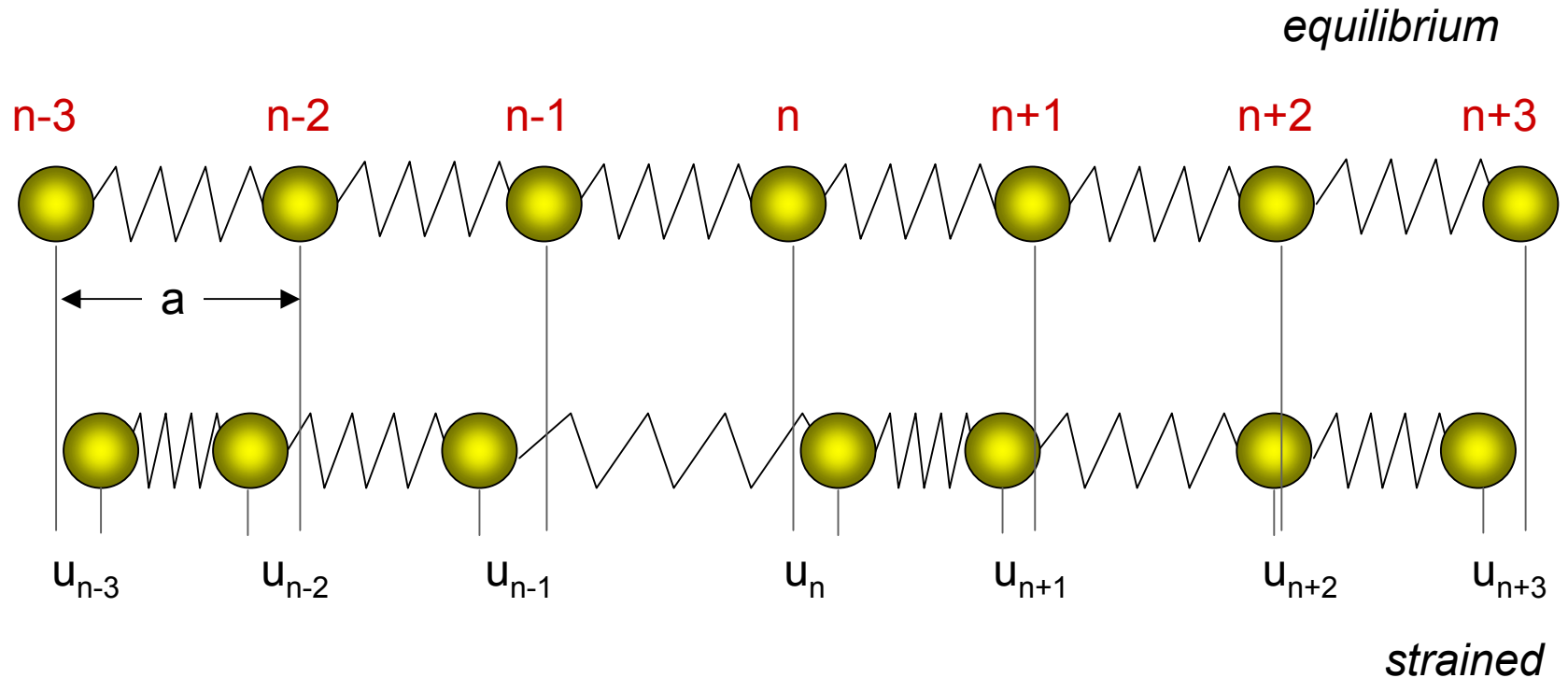
$$M\omega^2 = \sum_{m=-\infty}^{\infty} \tilde{D}(n, m) e^{ika(m-n)}$$

$$= \underbrace{\sum_{p=-\infty}^{\infty} \tilde{D}(p) e^{-ikap}}_{\text{Dynamical Matrix } D(k)}$$

$$\omega = \sqrt{\frac{D(k)}{M}}$$

Strain in a Discrete Lattice

Example of Nearest Neighbor Interactions



$$V = \sum_{p=-\infty}^{\infty} \frac{\alpha}{2} (u[p+1] - u[p])^2$$

Strain in a Discrete Lattice

Example of Nearest Neighbor Interactions

$$V = \sum_{p=-\infty}^{\infty} \frac{\alpha}{2} (u[p+1] - u[p])^2$$

$$\begin{aligned} \widetilde{D}(n, m) &= \left(\frac{\partial^2 V}{\partial u[n, t] \partial u[m, t]} \right)_{\text{eq}} = \widetilde{D}(n - m) \\ &= \frac{\partial}{\partial u[n, t]} \alpha (u[m] - u[m-1] - u[m+1] + u[m]) \end{aligned}$$

$$\widetilde{D}(0) = 2\alpha \quad \text{and} \quad \widetilde{D}(\pm 1) = -\alpha$$

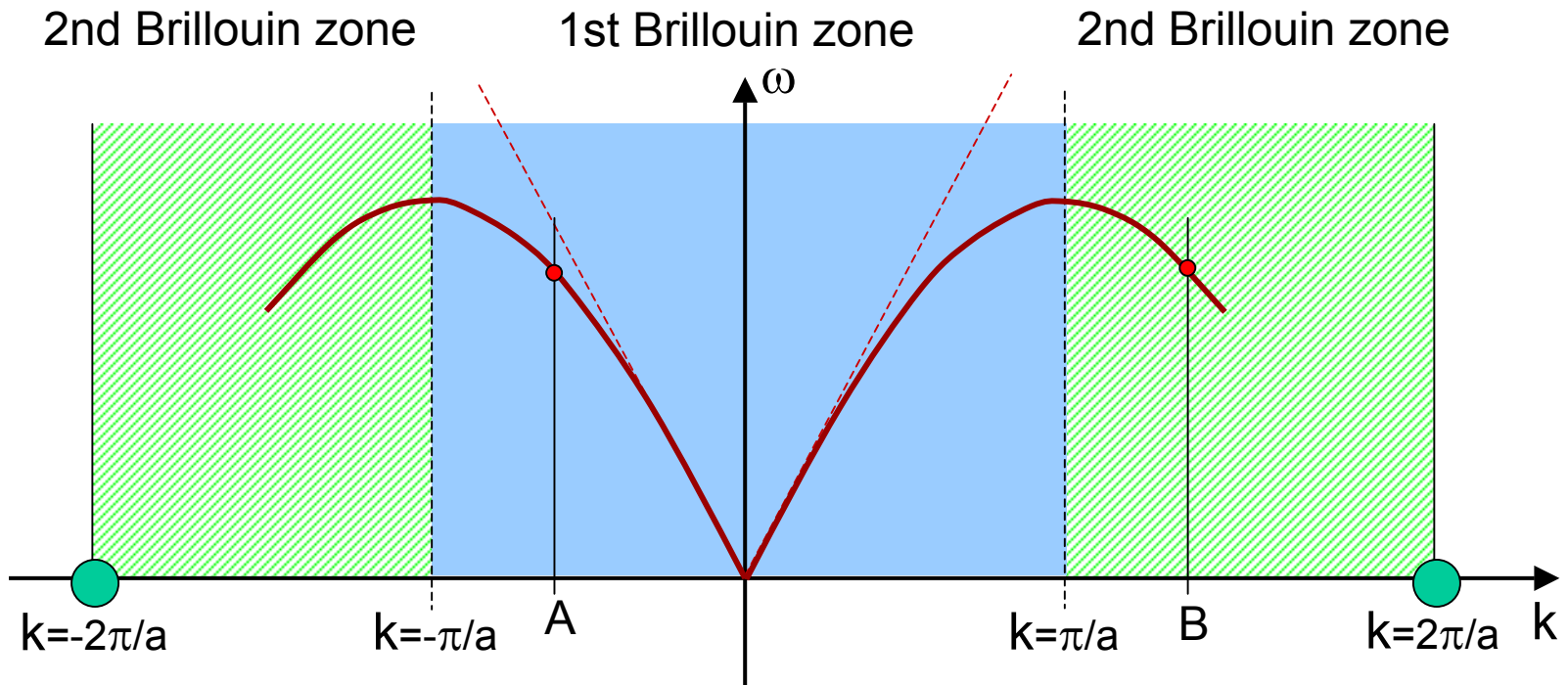
$$D(k) = \sum_{p=-\infty}^{\infty} \widetilde{D}(p) e^{-ikap}$$

$$D(k) = 2\alpha - \alpha e^{-ika} - \alpha e^{ika} = 2\alpha(1 - \cos ka) = 4\alpha \sin^2\left(\frac{ka}{2}\right)$$

Strain in a Discrete Lattice

Example of Nearest Neighbor Interactions

$$M\omega^2 = D(k) = 4\alpha \sin^2\left(\frac{ka}{2}\right) \quad \longrightarrow \quad \omega = 2\sqrt{\frac{\alpha}{M}} \left| \sin\left(\frac{ka}{2}\right) \right|$$



From what we know about Brillouin zones the points A and B (related by a reciprocal lattice vector) must be identical

$$\omega(k) = \omega(k + n2\pi/a)$$

Summary of Phonon Dispersion Calculation

- Taylor series expansion for total potential stored in all bonds
 - Neglect first order since in equilibrium $F=0$
 - Truncate expansion at second order, assume small amplitudes
- Determine harmonic matrix from potential energy
 - Represents bond stiffness

$$\tilde{D}(n, m) = \left(\frac{\partial^2 V}{\partial u[n, t] \partial u[m, t]} \right)_{\text{eq}}$$

- Assume time harmonic and discrete 'plane wave' solutions
- Determine dynamical matrix from harmonic matrix plus phase progression

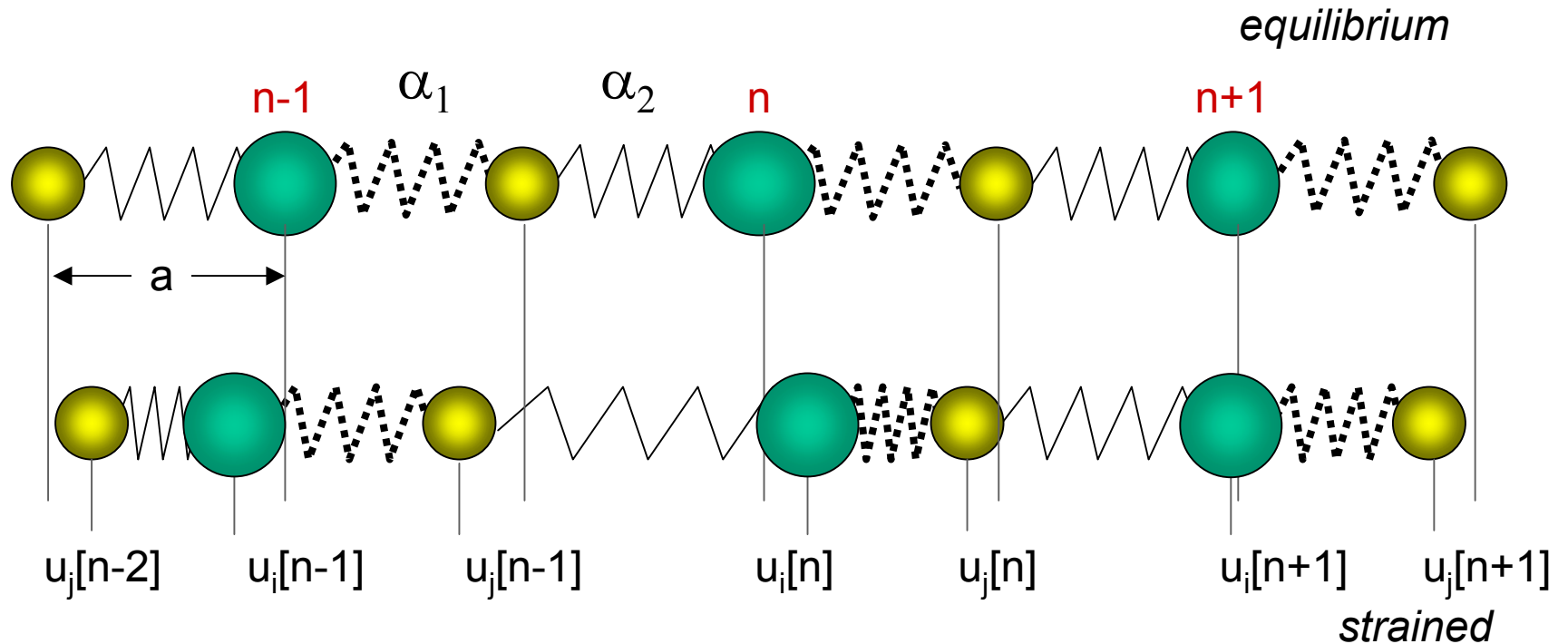
$$D(k) = \sum_{p=-\infty}^{\infty} \tilde{D}(p) e^{-ikap}$$

- Determine dispersion relation

$$\omega = \sqrt{\frac{D(k)}{M}}$$

Strain in a Discrete Lattice with Basis

Example of Nearest Neighbor Interactions



$$\begin{aligned}
 V(\{u[s, t]\}) = & V_0 + \sum_{i=1}^2 \sum_{m=-\infty}^{\infty} \left(\frac{\partial V}{\partial u_i[m, t]} \right)_{\text{eq}} u_i[m, t] \\
 & + \frac{1}{2} \sum_{i=1}^2 \sum_{j=1}^2 \sum_{p=-\infty}^{\infty} \sum_{m=-\infty}^{\infty} u_i[p, t] \left(\frac{\partial^2 V}{\partial u_i[p, t] \partial u_j[m, t]} \right)_{\text{eq}} u_j[m, t] + \dots
 \end{aligned}$$

Harmonic Matrix for 1-D Lattice with Basis

$$V(\{u[s, t]\}) = V_o + \frac{1}{2} \sum_{i=1}^2 \sum_{j=1}^2 \sum_{p=-\infty}^{\infty} \sum_{m=-\infty}^{\infty} u_i[p, t] \left(\frac{\partial^2 V}{\partial u_i[p, t] \partial u_j[m, t]} \right)_{\text{eq}} u_j[m, t]$$

$$\tilde{\mathbf{D}}_{i,j}(p, m) = \left(\frac{\partial^2 V}{\partial u_i[p, t] \partial u_j[m, t]} \right)_{\text{eq}}$$

$$V(\{u[s, t]\}) = V_o + \frac{1}{2} \sum_{i=1}^2 \sum_{j=1}^2 \sum_{p=-\infty}^{\infty} \sum_{m=-\infty}^{\infty} u_i[p, t] \tilde{\mathbf{D}}_{i,j}(p, m) u_j[m, t]$$

Equations of Motion

$$V(\{u[s, t]\}) = V_o + \frac{1}{2} \sum_{i=1}^2 \sum_{j=1}^2 \sum_{p=-\infty}^{\infty} \sum_{m=-\infty}^{\infty} u_i[p, t] \tilde{\mathbf{D}}_{i,j}(p, m) u_j[m, t]$$

The force on the ℓ^{th} basis atom in the n^{th} unit cell...

$$F_\ell[n, t] = -\frac{\partial V}{\partial u_\ell[n, t]}$$

$$M_\ell \frac{d^2}{dt^2} u_\ell[n] = -\frac{\partial}{\partial u_\ell[n]} V(\{u_i[s]\})$$

$$M_i \frac{d^2}{dt^2} u_i[n] = -\sum_{j=1}^2 \sum_{m=-\infty}^{\infty} \tilde{\mathbf{D}}_{i,j}(n, m) u_j[m]$$

$$\tilde{u}_i[n, t] = U_i[n, \omega] e^{-i\omega t}$$

Matrix Representation of Equations of Motion

$$M_i \omega^2 U_i[n] = \sum_{j=1}^2 \sum_{m=-\infty}^{\infty} \widetilde{\mathbf{D}}_{i,j}(n, m) U_j[m]$$

Can collect system of equations for each atom in the basis as a matrix...

$$\mathbf{U}[n] = \begin{pmatrix} U_1[n] \\ U_2[n] \end{pmatrix} \quad \mathbf{M} = \begin{pmatrix} M_1 & 0 \\ 0 & M_2 \end{pmatrix}$$

$$\omega^2 \mathbf{M} \mathbf{U}[n] = \sum_{\mathbf{m}=-\infty}^{\infty} \widetilde{\mathbf{D}}(\mathbf{n}, \mathbf{m}) \mathbf{U}[\mathbf{m}]$$

Plane Wave Solutions & the Dynamical Matrix

$$U[n + 1] = e^{ika}U[n]$$

$$U[n] = e^{ikna}U[0] = e^{ikna}\tilde{\epsilon}$$

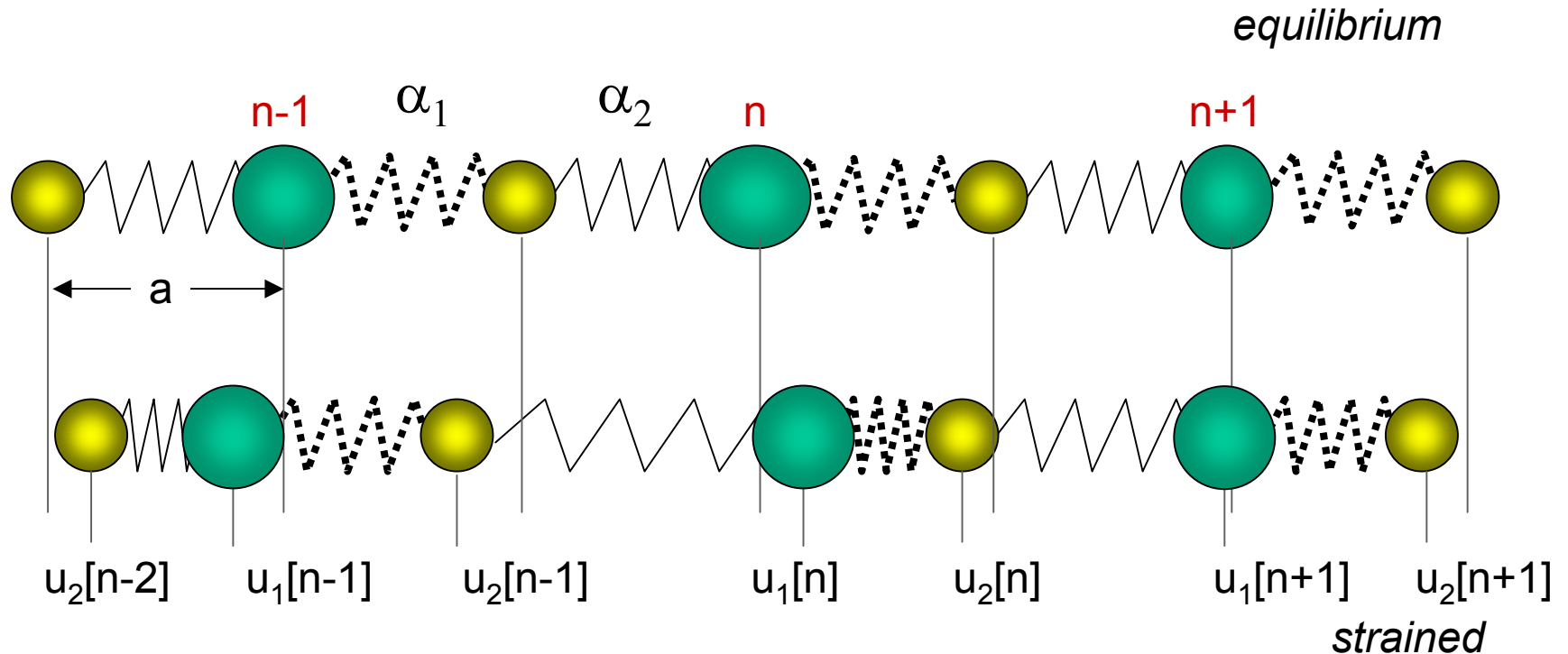
$$\omega^2 M\tilde{\epsilon} = D(k)\tilde{\epsilon}$$

$$D(k) = \sum_{m=-\infty}^{\infty} \tilde{D}(n-m)e^{ika(m-n)} = \sum_{p=-\infty}^{\infty} \tilde{D}(p)e^{-ikpa}$$

$$\left(M^{-1}D(k)\right)\tilde{\epsilon} = \omega^2\tilde{\epsilon}$$

Strain in a Discrete Lattice

Example of Nearest Neighbor Interactions



$$V = \dots + \frac{\alpha_1}{2} (u_1[s] - u_2[s])^2 + \frac{\alpha_2}{2} (u_1[s] - u_2[s-1])^2$$

$$+ \frac{\alpha_2}{2} (u_1[s+1] - u_2[s])^2 + \dots$$

Dynamical Matrix for 1-D Lattice with Basis

Example of Nearest Neighbor Coupling

$$V = \dots + \frac{\alpha_1}{2} (u_1[s] - u_2[s])^2 + \frac{\alpha_2}{2} (u_1[s] - u_2[s-1])^2$$

$$+ \frac{\alpha_2}{2} (u_1[s+1] - u_2[s])^2 + \dots$$

$$D_{i,j}(\mathbf{k}) = \sum_{\mathbf{R}_p} \left(\frac{\partial^2 V}{\partial u_i[\mathbf{R}_s + \mathbf{R}_p, t] \partial u_j[\mathbf{R}_s, t]} \right)_{\text{eq}} e^{-i\mathbf{k} \cdot \mathbf{R}_p}$$

$$D(\mathbf{k}) = \begin{matrix} & u_1 & & u_2 \\ u_1 & \left(\begin{array}{cc} \alpha_1 + \alpha_2 & -\alpha_1 - \alpha_2 e^{-ika} \\ -\alpha_1 - \alpha_2 e^{ika} & \alpha_1 + \alpha_2 \end{array} \right) & & \\ u_2 & & & \end{matrix}$$

Dispersion Relation for 1-D Lattice with Basis

Example of Nearest Neighbor Coupling

$$(\mathbf{M}^{-1}\mathbf{D}(\mathbf{k})) \vec{\epsilon} = \omega^2 \vec{\epsilon}$$

$$\mathbf{M}^{-1} = \begin{pmatrix} \frac{1}{M_1} & 0 \\ 0 & \frac{1}{M_2} \end{pmatrix} \quad \mathbf{D}(\mathbf{k}) = \begin{pmatrix} \alpha_1 + \alpha_2 & -\alpha_1 - \alpha_2 e^{-ika} \\ -\alpha_1 - \alpha_2 e^{ika} & \alpha_1 + \alpha_2 \end{pmatrix}$$

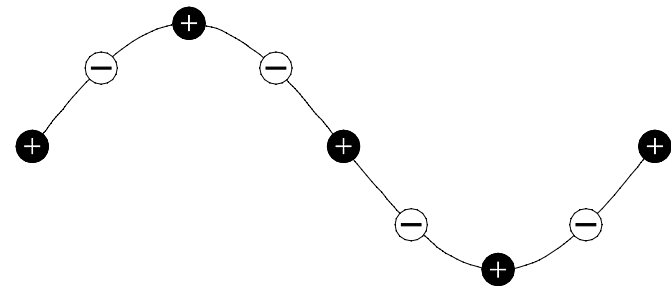
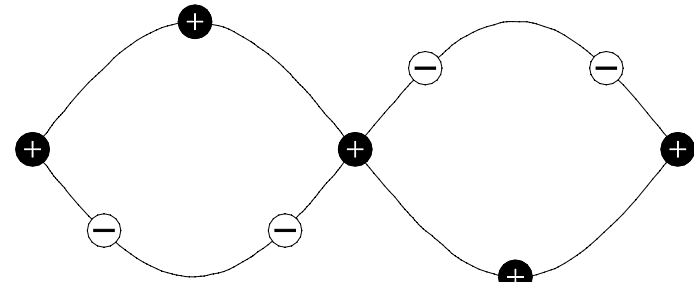
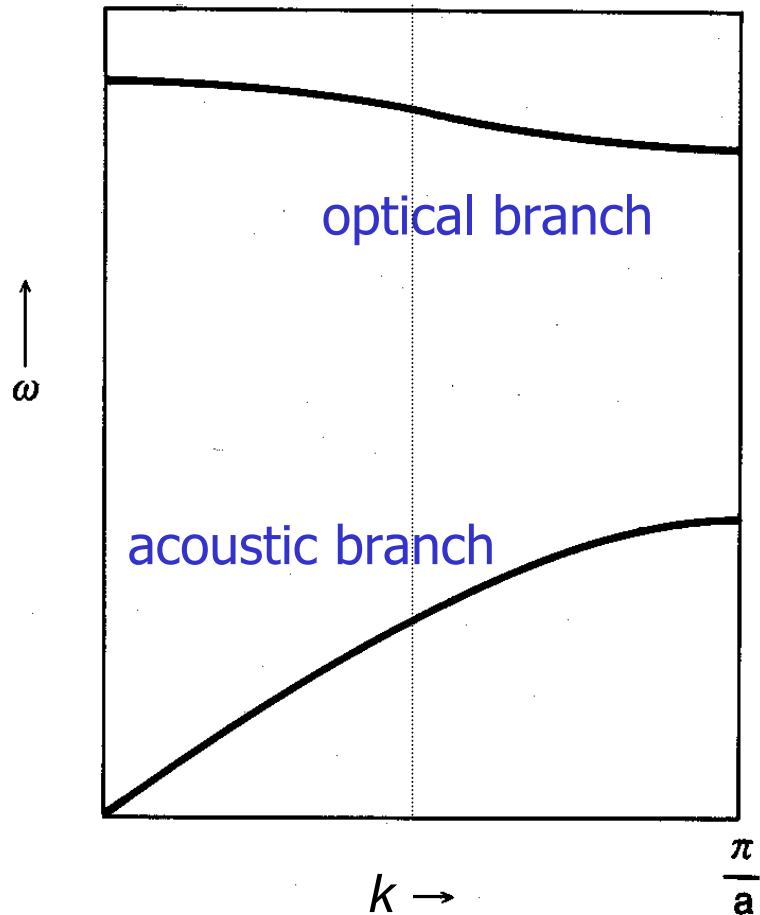
$$\begin{pmatrix} \frac{\alpha_1 + \alpha_2}{M_1} & -\frac{\alpha_1 + \alpha_2 e^{-ika}}{M_1} \\ -\frac{\alpha_1 + \alpha_2 e^{ika}}{M_2} & \frac{\alpha_1 + \alpha_2}{M_2} \end{pmatrix} \begin{pmatrix} \epsilon_1 \\ \epsilon_2 \end{pmatrix} = \omega^2 \begin{pmatrix} \epsilon_1 \\ \epsilon_2 \end{pmatrix}$$

$$\omega^2 = \frac{\alpha_1 + \alpha_2}{2} \left(\frac{1}{M_1} + \frac{1}{M_2} \right) \pm \left\{ \frac{(\alpha_1 + \alpha_2)^2 \left(\frac{1}{M_1} + \frac{1}{M_2} \right)^2}{4} - \frac{2\alpha_1\alpha_2(1 - \cos ka)}{M_1 M_2} \right\}^{1/2}$$

Dispersion Relation for 1-D Lattice with Basis

Example of Nearest Neighbor Coupling

$$\omega^2 = \frac{\alpha_1 + \alpha_2}{2} \left(\frac{1}{M_1} + \frac{1}{M_2} \right) \pm \left\{ \frac{(\alpha_1 + \alpha_2)^2 \left(\frac{1}{M_1} + \frac{1}{M_2} \right)^2}{4} - \frac{2\alpha_1\alpha_2(1 - \cos ka)}{M_1M_2} \right\}^{1/2}$$



Lattice Waves at $k=0$

Example of Nearest Neighbor Coupling

$$\begin{pmatrix} \frac{\alpha_1 + \alpha_2}{M_1} & -\frac{\alpha_1 + \alpha_2}{M_1} \\ -\frac{\alpha_1 + \alpha_2}{M_2} & \frac{\alpha_1 + \alpha_2}{M_2} \end{pmatrix} \begin{pmatrix} \epsilon_1(0) \\ \epsilon_2(0) \end{pmatrix} = \omega_\sigma^2 \begin{pmatrix} \epsilon_1(0) \\ \epsilon_2(0) \end{pmatrix}$$

$$\omega_1 = 0$$

$$\omega_2 = \sqrt{(\alpha_1 + \alpha_2) \left(\frac{1}{M_1} + \frac{1}{M_2} \right)}$$

$$\begin{pmatrix} \epsilon_1^{(1)}(0) \\ \epsilon_2^{(1)}(0) \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} \quad \begin{pmatrix} \epsilon_1^{(2)}(0) \\ \epsilon_2^{(2)}(0) \end{pmatrix} = \frac{1}{\sqrt{1 + (M_2/M_1)^2}} \begin{pmatrix} M_2/M_1 \\ -1 \end{pmatrix}$$

Lattice Waves at Small k

Example of Nearest Neighbor Coupling

$$\omega_1 = \left(\frac{\alpha_1 \alpha_2 a^2}{(\alpha_1 + \alpha_2)(M_1 + M_2)} \right)^{1/2} k \quad \omega_2 \approx \sqrt{(\alpha_1 + \alpha_2) \left(\frac{1}{M_1} + \frac{1}{M_2} \right)}$$

$$c_s = \left(\frac{\alpha_1 \alpha_2}{(\alpha_1 + \alpha_2)(M_1 + M_2)} \right)^{1/2}$$

$$\begin{pmatrix} \epsilon_1^{(1)}(0) \\ \epsilon_2^{(1)}(0) \end{pmatrix} \approx \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} \quad \text{for} \quad \omega_1 \approx c_s k$$

$$U_1[n+1] = e^{ika} U_1[n]$$

$$U_2[n+1] = e^{ika} U_2[n]$$

Lattice Waves Near Zone Boundary

Example of Nearest Neighbor Coupling

$$\begin{pmatrix} \frac{\alpha_1 + \alpha_2}{M} & \frac{\alpha_2 - \alpha_1}{M} \\ \frac{\alpha_2 - \alpha_1}{M} & \frac{\alpha_1 + \alpha_2}{M} \end{pmatrix} \begin{pmatrix} \epsilon_1^{(i)} \left(\frac{\pi}{a}\right) \\ \epsilon_2^{(i)} \left(\frac{\pi}{a}\right) \end{pmatrix} = \omega_i^2 \begin{pmatrix} \epsilon_1^{(i)} \left(\frac{\pi}{a}\right) \\ \epsilon_2^{(i)} \left(\frac{\pi}{a}\right) \end{pmatrix}$$

$$\omega_1 = \sqrt{\frac{2\alpha_1}{M}} \quad \text{and} \quad \omega_2 = \sqrt{\frac{2\alpha_2}{M}}$$

$$\begin{pmatrix} \epsilon_1^{(1)} \left(\frac{\pi}{a}\right) \\ \epsilon_2^{(1)} \left(\frac{\pi}{a}\right) \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix} \quad \text{for} \quad \omega_1^2 = \frac{2\alpha_1}{M}$$

$$\begin{pmatrix} \epsilon_1^{(2)} \left(\frac{\pi}{a}\right) \\ \epsilon_2^{(2)} \left(\frac{\pi}{a}\right) \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} \quad \text{for} \quad \omega_2^2 = \frac{2\alpha_2}{M}$$

Dispersion Relation for 3-D Lattices

