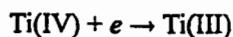


Model Solutions to 3.53 Problem Set 7

(sample problems; not to be submitted as an assignment)

Problem 10.7 For



the following experimental conditions are given.

$$n = 1$$

$$C_{\text{Ti(IV)}}^* = 3.36 \text{ mM}$$

$$T = 25^\circ\text{C}$$

$$E_{dc} = E_{1/2} = -0.290 \text{ V vs SCE}$$

$$D_O = D_R = 6.6 \times 10^{-6} \text{ cm}^2/\text{s}$$

From equation (10.5.25),

$$\begin{aligned} [\cot \phi]_{E_{1/2}} &= 1 + \left(\frac{D_O^\beta D_R^\alpha}{2} \right)^{1/2} \frac{\omega^{1/2}}{k^0} & (10.5.25) \\ &= 1 + \sqrt{\frac{D}{2}} \frac{\omega^{1/2}}{k^0} & (1) \end{aligned}$$

since $\alpha + \beta = 1$ (footnote 4, page 95; discussion after equation (10.5.25), page 395). A plot of $[\cot \phi]_{E_{1/2}}$ vs $\omega^{1/2}$ will have, from equation (1),

$$\text{slope} = \sqrt{\frac{D}{2}} \frac{1}{k^0} \quad (2)$$

From Figure 10.5.5.

$$\text{slope} \approx \frac{4.45}{88 \text{ s}^{-1/2}} = 5.06 \times 10^{-2} \text{ s}^{1/2} = \sqrt{\frac{D}{2}} \frac{1}{k^0} \quad (3)$$

Solving equation (3) for k^0 leads to a value of $k^0 = 3.6 \times 10^{-2} \text{ cm/s}$. From equation (10.5.24), at $[\cot \phi]_{\text{max}}$,

$$E_{dc} - E_{1/2} = \frac{RT}{F} \ln \frac{\alpha}{\beta} = (0.0257 \text{ V}) \ln \frac{\alpha}{\beta} \approx -0.016 \text{ V at } [\cot \phi]_{\text{max}} \quad (4)$$

Substituting $(1 - \alpha) = \beta$ in equation (4) and algebraic manipulation leads to

$$\frac{\alpha}{1 - \alpha} = \exp[-0.6226] = 0.54 \quad (5)$$

which solves to $\alpha = 0.35$.

Problem 10.8 From equation (10.3.9),

$$\phi = \tan^{-1} \left[\frac{\sigma/\omega^{1/2}}{R_{ct} + \sigma/\omega^{1/2}} \right] \quad (10.3.9)$$

where

$$\sigma = \frac{RT}{n^2 F^2 A \sqrt{2}} \left[\frac{1}{\sqrt{D_O} C_O^*} + \frac{1}{\sqrt{D_R} C_R^*} \right] \quad (10.3.10)$$

$$R_{ct} = \frac{RT}{F i_0} \quad (10.3.2)$$

$$i_0 = n F A k^0 C_O^{*(1-\alpha)} C_R^{*\alpha} \quad (3.4.6)$$

It is given that $k^0 = 2.2 \pm 0.3 \text{ cm/s}$, $\alpha = 0.70$, $D_O = 1.02 \times 10^{-5} \text{ cm}^2/\text{s}$, $n = 1$, and $T = 295 \pm 2 \text{ K}$. For $n = 1$, substitution of equations (10.3.10), (10.3.2), and (3.4.6) into equation (10.3.9) yields

$$\begin{aligned} \phi &= \tan^{-1} \left[\frac{\frac{RT}{n^2 F^2 A \sqrt{2\omega}} \left[\frac{1}{\sqrt{D_O} C_O^*} + \frac{1}{\sqrt{D_R} C_R^*} \right]}{\frac{RT}{F n F A k^0 C_O^{*(1-\alpha)} C_R^{*\alpha}} + \frac{RT}{n^2 F^2 A \sqrt{2\omega}} \left[\frac{1}{\sqrt{D_O} C_O^*} + \frac{1}{\sqrt{D_R} C_R^*} \right]} \right] \quad (1) \\ &= \tan^{-1} \left[\frac{k^0 C_O^{*(1-\alpha)} C_R^{*\alpha} \left[\frac{1}{\sqrt{D_O} C_O^*} + \frac{1}{\sqrt{D_R} C_R^*} \right]}{\sqrt{2\omega} + k^0 C_O^{*(1-\alpha)} C_R^{*\alpha} \left[\frac{1}{\sqrt{D_O} C_O^*} + \frac{1}{\sqrt{D_R} C_R^*} \right]} \right] \end{aligned}$$

Let $C_O^* = C_R^* = C^*$ and $D_R = D_O$.

$$\phi = \tan^{-1} \left[\frac{\frac{2k^0}{\sqrt{D_O}}}{\sqrt{2\omega} + \frac{2k^0}{\sqrt{D_O}}} \right] \quad (2)$$

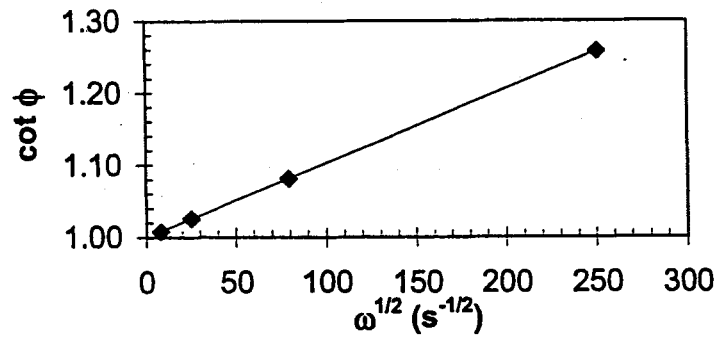
It is given that $k^0 = 2.2 \pm 0.3 \text{ cm/s}$, $\alpha = 0.70$, $D_O = 1.02 \times 10^{-5} \text{ cm}^2/\text{s}$, and $T = 295 \pm 2 \text{ K}$, such that $k^0/\sqrt{D_O} = 688 \text{ s}^{-1/2}$. For several decades of ω , ϕ is tabulated below.

$\omega/2\pi$	ω	$\phi(\text{rad})$	$\phi(\text{deg})$	$\omega^{1/2}$	$\cot \phi$
10	62.8	0.7813	44.77	7.93	1.008
100	628	0.7727	44.27	25.07	1.026
1000	6283	0.7463	42.76	79.27	1.081
10000	62831	0.6718	38.49	250.66	1.258

For reversible reactions, $\phi = 45^\circ$. For $k^0 = 2.2 \pm 0.3 \text{ cm/s}$, the reaction will be reversible at low frequencies, as is consistent with the data in the table where $\phi \rightarrow 45^\circ$ as ω decreases.

A plot of $\cot \phi = 1/\tan \phi$ versus $\omega^{1/2}$ is shown. Note that $E = E_{1/2} = E^0$ when $D_O = D_R$; then, k^0 is the operative heterogeneous rate. For these conditions, equation (10.5.25) applies, and it simplifies as shown for $D = D_O = D_R$ where $\beta = 1 - \alpha$.

$$\begin{aligned}
[\cot \phi]_{E_{1/2}} &= 1 + \left[\frac{D_O^\beta D_R^\alpha}{2} \right]^{1/2} \frac{\omega^{1/2}}{k^0} \\
&= 1 + \frac{D^{1/2}}{\sqrt{2}k^0} \omega^{1/2}
\end{aligned}
\tag{10.5.25}$$



Regression yields $\cot \phi = 1.03 \times 10^{-3} \omega^{1/2} + 1.0000$. The slope = $\sqrt{D/2}/k^0$; for the values here, $\sqrt{D/2}/k^0 = 1.03 \times 10^{-3} \text{ s}^{1/2}$.

Consider Figure 10.3.3, which shows the real and imaginary vectors that define the response for a quasireversible electron transfer. The real vector, measured along the same vector as \dot{E}_{ac} for a phase angle of 0° , is $R_{ct} + \sigma/\omega^{1/2}$. The vector 90° out of phase defines the imaginary term, $\sigma/\omega^{1/2}$. The ratio of these two terms defines $\cot \phi$. From equation (10.3.9).

$$\cot \phi = \omega R_s C_s = \frac{R_{ct} + \sigma/\omega^{1/2}}{\sigma/\omega^{1/2}}
\tag{3}$$

Thus, the ratio of a current measurement on the real axis made at 0° displacement with respect to \dot{E}_{ac} and a second current measurement 90° out of phase (quadrature current) will yield $\cot \phi$. Note that this assumes effects from uncompensated solution resistance and double layer charging are negligible.

To make a good measurement of k^0 , the frequency must be high enough that the measured value of ϕ must be less than 45° . As above, this condition is favored by higher frequency (faster measurements). Here, frequencies greater than 10 kHz are needed to reduce ϕ by at least one degree. Commercial instrumentation is available that generate frequencies of 20 MHz.