MATH 18.152 COURSE NOTES - CLASS MEETING # 9

18.152 Introduction to PDEs, Fall 2011	Professor: Jared Speck

Class Meeting # 9: Poisson's Formula, Harnack's Inequality, and Liouville's Theorem

1. Representation Formula for Solutions to Poisson's Equation

We now derive our main representation formula for solution's to Poisson's equation on a domain Ω .

Theorem 1.1 (Representation formula for solutions to the boundary value Poisson equation). Let Ω be a domain with a smooth boundary, and assume that $f \in C^2(\overline{\Omega})$ and $g \in C(\partial\Omega)$. Then the unique solution $u \in C^2(\Omega) \cap C(\overline{\Omega})$ to

(1.0.1)
$$\Delta u(x) = f(x), \qquad x \in \Omega \subset \mathbb{R}^n,$$
$$u(x) = g(x), \qquad x \in \partial \Omega.$$

can be represented as

(1.0.2)
$$u(x) = \int_{\Omega} f(y)G(x,y) d^{n}y + \int_{\partial\Omega} g(\sigma) \underbrace{\nabla_{\hat{N}(\sigma)}G(x,\sigma)}_{Poisson \ kernel} d\sigma,$$

where G(x, y) is the Green function for Ω .

Proof. Applying the **Representation formula for** u Proposition, we have that

$$(1.0.3) \qquad u(x) = \int_{\Omega} \Phi(x-y)f(y) d^{n}y - \int_{\partial\Omega} \Phi(x-\sigma)\nabla_{\hat{N}(\sigma)}u(\sigma) d\sigma + \int_{\partial\Omega} g(\sigma)\nabla_{\hat{N}(\sigma)}\Phi(x-\sigma) d\sigma.$$

Recall also that

(1.0.4)
$$G(x,y) = \Phi(x-y) - \phi(x,y),$$

where

(1.0.5)
$$\Delta_y \phi(x, y) = 0, \qquad x \in \Omega,$$

and

(1.0.6)
$$G(x, \sigma) = 0$$
 when $x \in \Omega$ and $\sigma \in \partial \Omega$.

The expression (1.0.3) is not very useful since don't know the value of $\nabla_{\hat{N}(\sigma)} u(\sigma)$ along $\partial\Omega$. To fix this, we will use Green's identity. Applying Green's identity to the functions u(y) and $\phi(x, y)$, and recalling that $\Delta_y \phi(x, y) = 0$ for each fixed $x \in \Omega$, we have that

(1.0.7)
$$0 = \int_{\Omega} \phi(x,y) \underbrace{\widetilde{f(y)}}_{\partial u} d^{n}y - \int_{\partial \Omega} \underbrace{\phi(x,\sigma)}_{\partial v} \nabla_{\hat{N}} u(\sigma) d\sigma + \int_{\partial \Omega} \underbrace{g(\sigma)}_{\partial v} \nabla_{\hat{N}} \phi(x,\sigma) d\sigma.$$

Subtracting (1.0.7) from (1.0.3), and using (1.0.4), we deduce the formula (1.0.2).

2. Poisson's Formula

Let's compute the Green function G(x, y) and Poisson kernel $P(x, \sigma) \stackrel{\text{def}}{=} \nabla_{\hat{N}} G(x, \sigma)$ from (1.0.2) in the case that $\Omega \stackrel{\text{def}}{=} B_R(0) \subset \mathbb{R}^3$ is a ball of radius R centered at the origin. We'll use a technique called the *method of images* that works for special domains.

Warning 2.0.1. Brace yourself for a bunch of tedious computations that at the end of the day will lead to a very nice expression.

The basic idea is to hope that $\phi(x, y)$ from the decomposition $G(x, y) = \Phi(x - y) - \phi(x, y)$, where $\phi(x, y)$ is viewed as a function of x that depends on the parameter y, is equal to the Newtonian potential generated by some "imaginary charge" q placed at a point $x^* \in B_R^c(0)$. To ensure that $G(x, \sigma) = 0$ when $\sigma \in \partial B_R(0)$, q and x^* have to be chosen so that along the boundary $\{y \in \mathbb{R}^3 \mid |y| = R\}$, $\phi(x, y) = \frac{1}{4\pi |x-y|}$. In a nutshell, we guess that

(2.0.8)
$$G(x,y) = -\frac{1}{4\pi |x-y|} + \underbrace{\frac{q}{4\pi |x^*-y|}}_{\phi(x,y)?},$$

and we try to solve for q and x^* so that G(x, y) vanishes when |y| = R.

Remark 2.0.1. Note that $\Delta_y \frac{q}{4\pi |x^*-y|} = 0$, which is one of the conditions necessary for constructing G(x, y).

By the definition of G(x, y), we must have G(x, y) = 0 when |y| = R, which implies that

(2.0.9)
$$\frac{1}{4\pi|x-y|} = \frac{q}{4\pi|x^*-y|}.$$

Simple algebra then leads to

(2.0.10)
$$|x^* - y|^2 = q^2 |x - y|^2.$$

When |y| = R, we use (2.0.10) to compute that

(2.0.11)
$$|x^*|^2 - 2x^* \cdot y + R^2 = |x^* - y|^2 = q^2 |x - y|^2 = q^2 (|x|^2 - 2x \cdot y + R^2),$$

where \cdot denotes the Euclidean dot product. Then performing simple algebra, it follows from (2.0.11) that

(2.0.12)
$$|x^*|^2 + R^2 - q^2(R^2 + |x|^2) = 2y \cdot (x^* - q^2x).$$

Now since the left-hand side of (2.0.12) does not depend on y, it must be the case that the right-hand side is always 0. This implies that $x^* = q^2 x$, and also leads to the equation

(2.0.13)
$$q^4|x|^2 - q^2(R^2 + |x|^2) + R^2 = 0.$$

Solving (2.0.13) for q, we finally have that

$$(2.0.14) q = \frac{R}{|x|},$$

(2.0.15)
$$x^* = \frac{R^2}{|x|^2}x.$$

Therefore,

(2.0.16)
$$\phi(x,y) = \frac{1}{4\pi} \frac{R}{|x| |\frac{R^2}{|x|^2} x - y|},$$

(2.0.17)
$$\phi(0,y) = \frac{1}{4\pi R},$$

where we took a limit as $x \to 0$ in (2.0.16) to derive (2.0.17). Next, using (2.0.8), we have

(2.0.18)
$$G(x,y) = -\frac{1}{4\pi |x-y|} + \frac{1}{4\pi} \frac{R}{|x| |\frac{R^2}{|x|^2} x - y|}, \qquad x \neq 0,$$

(2.0.19)
$$G(0,y) = -\frac{1}{4\pi |y|} + \frac{1}{4\pi R}.$$

For future use, we also compute that

(2.0.20)
$$\nabla_y G(x,y) = -\frac{x-y}{4\pi |x-y|^3} + \frac{1}{4\pi} \frac{R}{|x|} \frac{x^*-y}{|x^*-y|^3}$$

Now when $\sigma \in \partial B_R(0)$, (2.0.10) and (2.0.14) imply that

(2.0.21)
$$|x^* - \sigma| = \frac{R}{|x|}|x - \sigma|.$$

Therefore, using (2.0.20) and (2.0.21), we compute that

$$(2.0.22) \qquad \nabla_{\sigma}G(x,\sigma) = -\frac{x-\sigma}{4\pi|x-\sigma|^3} + \frac{1}{4\pi}\frac{|x|^2}{R^2}\frac{x^*-\sigma}{|x-\sigma|^3} = -\frac{x-\sigma}{4\pi|x-\sigma|^3} + \frac{1}{4\pi}\frac{|x|^2}{R^2}\frac{\frac{R^2}{|x|^2}x-\sigma}{|x-\sigma|^3} \\ = \frac{\sigma}{4\pi|x-\sigma|^3}\Big(1-\frac{|x|^2}{R^2}\Big).$$

Using (2.0.22) and the fact that $\hat{N}(\sigma) = \frac{1}{R}\sigma$, we deduce

(2.0.23)
$$\nabla_{\hat{N}(\sigma)} G(x,\sigma) \stackrel{\text{def}}{=} \nabla_{\sigma} G(x,\sigma) \cdot \hat{N}(\sigma) = \frac{R^2 - |x|^2}{4\pi R} \frac{1}{|x-\sigma|^3}.$$

Remark 2.0.2. If the ball were centered at the point $p \in \mathbb{R}^3$ instead of the origin, then the formula (2.0.23) would be replaced with

(2.0.24)
$$\nabla_{\hat{N}(\sigma)} G(x,\sigma) \stackrel{\text{def}}{=} \nabla_{\sigma} G(x,\sigma) \cdot \hat{N}(\sigma) = -\frac{R^2 - |x-p|^2}{4\pi R} \frac{1}{|x-\sigma|^3}.$$

Let's summarize this by stating a lemma.

Lemma 2.0.1. The Green function for a ball $B_R(p) \subset \mathbb{R}^3$ is

(2.0.25a)

(2.0.25a)
$$G(x,y) = -\frac{1}{4\pi |x-y|} + \frac{1}{4\pi} \frac{R}{|x-p| \left| \frac{R^2}{|x-p|^2} (x-p) - (y-p) \right|},$$

(2.0.25b)
$$G(p,y) = -\frac{1}{4\pi |y-p|} + \frac{1}{4\pi R}.$$

Furthermore, if $x \in B_R(p)$ and $\sigma \in \partial B_R(p)$, then

(2.0.25c)
$$\nabla_{\hat{N}(\sigma)} G(x,\sigma) = \frac{R^2 - |x-p|^2}{4\pi R} \frac{1}{|x-\sigma|^3}.$$

We can now easily derive a representation formula for solutions to the Laplace equation on a ball.

 $x \neq p$,

Theorem 2.1 (Poisson's formula). Let $B_R(p) \subset \mathbb{R}^3$ be a ball of radius R centered at p = (p^1, p^2, p^3) , and let $x = (x^1, x^2, x^3)$ denote a point in \mathbb{R}^3 . Let $g \in C(\partial B_R(p))$. Then the unique solution $u \in C^2(B_R(p)) \cap C(\overline{B}_R(p))$ of the PDE

(2.0.26)
$$\begin{cases} \Delta u(x) = 0, & x \in B_R(p), \\ u(x) = g(x), & x \in \partial B_R(p), \end{cases}$$

can be represented using the Poisson formula:

(2.0.27)
$$u(x) = \frac{R^2 - |x - p|^2}{4\pi R} \int_{\partial B_R(p)} \frac{g(\sigma)}{|x - \sigma|^3} \, d\sigma.$$

Remark 2.0.3. In *n* dimensions, the formula (2.0.27) gets replaced with

(2.0.28)
$$u(x) = \frac{R^2 - |x-p|^2}{\omega_n R} \int_{\partial B_R(p)} \frac{g(\sigma)}{|x-\sigma|^n} d\sigma,$$

where as usual, ω_n is the surface area of the unit ball in \mathbb{R}^n .

Proof. The identity (2.0.27) follows immediately from Theorem 1.1 and Lemma 2.0.1.

3. HARNACK'S INEQUALITY

We will now use some of our tools to prove a famous inequality for Harmonic functions. The theorem provides some estimates that place limitations on how slow/fast harmonic functions are allowed to grow.

Theorem 3.1 (Harnack's inequality). Let $B_R(0) \subset \mathbb{R}^n$ be the ball of radius R centered at the origin, and let $u \in C^2(B_R(0)) \cap C(\overline{B}_R(0))$ be the unique solution to (2.0.26). Assume that u is non-negative on $\overline{B}_R(0)$. Then for any $x \in B_R(0)$, we have that

(3.0.29)
$$\frac{R^{n-2}(R-|x|)}{(R+|x|)^{n-1}}u(0) \le u(x) \le \frac{R^{n-2}(R+|x|)}{(R-|x|)^{n-1}}u(0).$$

Proof. We'll do the proof for n = 3. The basic idea is to combine the Poisson representation formula with simple inequalities and the mean value property. By Theorem 2.1, we have that

(3.0.30)
$$u(x) = \frac{R^2 - |x|^2}{4\pi R} \int_{\partial B_R(0)} \frac{g(\sigma)}{|x - \sigma|^3} d\sigma$$

By the triangle inequality, for $\sigma \in \partial B_R(0)$ (i.e. $|\sigma| = R$), we have that $|x| - R \le |x - \sigma| \le |x| + R$. Applying the first inequality to (3.0.30), and using the non-negativity of g, we deduce that

(3.0.31)
$$u(x) \le \frac{R+|x|}{R^2 - |x|^2} \frac{1}{4\pi R} \int_{\partial B_R(0)} g(\sigma) \, d\sigma.$$

Now recall that by the mean value property, we have that

(3.0.32)
$$u(0) = \frac{1}{4\pi R^2} \int_{\partial B_R(0)} g(\sigma) \, d\sigma.$$

Thus, combining (3.0.31) and (3.0.32), we have that

(3.0.33)
$$u(x) \le \frac{R(R+|x|)}{(R-|x|)^2}u(0),$$

which implies one of the inequalities in (3.0.29). The other one can be proved similarly using the remaining triangle inequality.

We now prove a famous consequence of Harnack's inequality. The statement is also often proved in introductory courses in complex analysis, and it plays a central role in some proofs of the fundamental theorem of algebra.

Corollary 3.0.2 (Liouville's theorem). Suppose that $u \in C^2(\mathbb{R}^n)$ is harmonic on \mathbb{R}^n . Assume that there exists a constant M such that $u(x) \ge M$ for all $x \in \mathbb{R}^n$, or such that $u(x) \le M$ for all $x \in \mathbb{R}^n$. Then u is a constant-valued function.

Proof. We first consider the case that $u(x) \ge M$. Let $v \stackrel{\text{def}}{=} u + |M|$. Observe that $v \ge 0$ is harmonic and verifies the hypotheses of Theorem 3.1. Thus, by (3.0.29), if $x \in \mathbb{R}^n$ and R is sufficiently large, we have that

(3.0.34)
$$\frac{R^{n-2}(R-|x|)}{(R+|x|)^{n-1}}v(0) \le v(x) \le \frac{R^{n-2}(R+|x|)}{(R-|x|)^{n-1}}v(0).$$

Allowing $R \to \infty$ in (3.0.34), we conclude that v(x) = v(0). Thus, v is a constant-valued function (and therefore u is too).

To handle the case $u(x) \leq M$, we simply consider the function $w(x) \stackrel{\text{def}}{=} -u(x) + |M|$ in place of v(x), and we argue as above.

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