

Lecture 24: Laplacian Growth II

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In the last lecture, we looked at the continuous dynamics of a freely moving domain. Using a mapping from the boundary in the z -plan to the unit circle in the ω -plan, $z = g(\omega, t)$, we have proved the following relation for the finite diffusive growth:

$$\operatorname{Re}(\overline{\omega g'} g_t) = 1 \quad \text{on } |\omega| = 1 \quad (1)$$

1 Laurent Expansion of $g(\omega, t)$

We know from the previous lecture that $g(\omega, t)$ must be a univalent (1-1, conformal) map from $\{|\omega| \geq 1\}$ to $\Omega_z(t)$. Thus, we write the Laurent Expansion:

$$g(\omega, t) = \sum_{n=-\infty}^1 A_n(t) \omega^n \quad \text{for } |\omega| \geq 1$$

where the sum goes up to $n = 1$ to ensure that g is a 1-1 map. We calculate:

$$\begin{aligned} g'(\omega, t) &= \sum_{n=-\infty}^1 n A_n(t) \omega^{n-1} \\ \omega g'(\omega, t) &= \sum_{n=-\infty}^1 n A_n(t) \omega^n \\ \overline{\omega g'(\omega, t)} &= \sum_{n=-\infty}^1 \overline{n A_n(t)} \omega^{-n} \end{aligned}$$

since $\bar{\omega} = \omega^{-1}$ for $|\omega| = 1$. Also:

$$g_t(\omega, t) = \sum_{n=-\infty}^1 \dot{A}_n(t) \omega^n$$

to finally write from (1):

$$\operatorname{Re} \left(\sum_{n, n'=-\infty}^1 \overline{n A_n(t)} \dot{A}_{n'}(t) \omega^{n'-n} \right) = 1 \quad (2)$$

1.1 The Area Equation

Since the Laurent coefficients are unique, one must verify the following relation from (2) for the terms $n = n'$:

$$\begin{aligned} & \operatorname{Re} \left(\sum_{n=-\infty}^1 n \overline{A_n(t)} \dot{A}_n(t) \right) = 1 \\ &= \frac{1}{2} \sum_{n=-\infty}^1 n \left(\overline{A_n(t)} \dot{A}_n(t) + A_n(t) \overline{\dot{A}_n(t)} \right) \\ &= \frac{1}{2} \frac{d}{dt} \sum_{n=-\infty}^1 n |A_n(t)|^2 \end{aligned}$$

that is:

$$\sum_{n=-\infty}^1 n |A_n(t)|^2 = 2t + \sum_{n=-\infty}^1 n |A_n(0)|^2 \quad (3)$$

Aside, we can notice that the area $\mathcal{A}(t)$ of the domain $\mathcal{C}_z(t)$ can be computed the following way:

$$\mathcal{A}(t) = \frac{1}{2i} \oint_{\mathcal{C}_z(t)} \bar{z} dz = \frac{1}{2i} \oint_{|\omega|=1} \bar{g} g' d\omega$$

by using the change of variable $z = g(\omega, t)$, with $\bar{z} = \bar{g}$ and $dz = g' d\omega$. Plugging the Laurent expansion of $g(\omega, t)$:

$$\begin{aligned} \frac{1}{2i} \oint_{|\omega|=1} \bar{g} g' d\omega &= \frac{1}{2i} \oint_{|\omega|=1} \left(\sum_{n,n'=-\infty}^1 n \overline{A_n(t)} A_{n'}(t) \bar{\omega}^n \omega^{n'-1} \right) d\omega \\ &= \frac{1}{2i} \sum_{n,n'=-\infty}^1 n \overline{A_n(t)} A_{n'}(t) \oint_{|\omega|=1} \frac{d\omega}{\omega^{n-n'+1}} \\ &= \pi \sum_{n=-\infty}^1 |A_n(t)|^2 \end{aligned}$$

since:

$$\oint_{|\omega|=1} \frac{d\omega}{\omega^{n-n'+1}} = \begin{cases} 2\pi i & \text{if } n = n' \\ 0 & \text{if } n \neq n' \end{cases}$$

Finally, (3) can be written:

$$\mathcal{A}(t) = \pi \sum_{n=-\infty}^1 n |A_n(t)|^2 = \mathcal{A}(0) + 2\pi t \quad (4)$$

which finds its physical interpretation from the fact that the total flux from ∞ is constant and equal to 2π .

1.2 Other terms

We are now interested by the terms of (2) for which $n \neq n'$:

$$\sum_{n=-\infty}^{1-k} \left[n \dot{A}_{n+k}(t) \overline{A_n(t)} + (n+k) A_{n+k}(t) \overline{\dot{A}_n(t)} \right] = 0 \quad \text{for } k = n' - n = 1, 2, \dots \quad (5)$$

Using Laurent expansion, we simplified considerably the problem, since we went from an integral partial differential equation (1) to a system of first order non-linear ordinary differential equation.

2 Exact Solutions

2.1 Circle

For a simple growing circle:

$$g(\omega, t) = A_1(t)\omega$$

From (4), we get:

$$\pi |A_1(t)|^2 = \pi + 2\pi t$$

Since $\mathcal{A}(0) = \pi$ for the unit circle. Choosing $A_1(t)$ real:

$$A_1(t) = \sqrt{1 + 2t}$$

is simply the radius of the growing circle.

2.2 Ellipse

For the ellipse, we verify:

$$g(\omega, t) = A_1(t)\omega + \frac{A_{-1}(t)}{\omega}$$

Indeed, this mapping is the sum of 2 circular, counter rotating mapping with same rotation angle. Assuming $|A_1(t)| < |A_{-1}(t)|$, this produces an ellipse with semi-major $|A_1(t)| + |A_{-1}(t)|$ and semi-minor $|A_1(t)| - |A_{-1}(t)|$. Without loss of generality, we will write:

$$A_1(0) = 1$$

$$A_{-1}(0) = c$$

where c is real and $0 < c < 1$. It can be shown that $A_1(t)$ and $A_{-1}(t)$ remain real at all times. The area equation (4) becomes:

$$|A_1(t)|^2 - |A_{-1}(t)|^2 = 2t + 1 - c^2$$

and (5) gives:

$$-\dot{A}_{-1}(t)A_1(t) + A_1(t)\dot{A}_{-1}(t) = 0 \quad \Rightarrow \quad \frac{\dot{A}_1(t)}{A_1(t)} = \frac{\dot{A}_{-1}(t)}{A_{-1}(t)}$$

and after integration:

$$A_1(t) = \frac{A_{-1}(t)}{c}$$

Finally, we write:

$$g(\omega, t) = \sqrt{1 + \frac{2t}{1-c^2}} \left(\omega + \frac{c}{\omega} \right)$$

Taking $c = 1$ leads to the circular mapping.

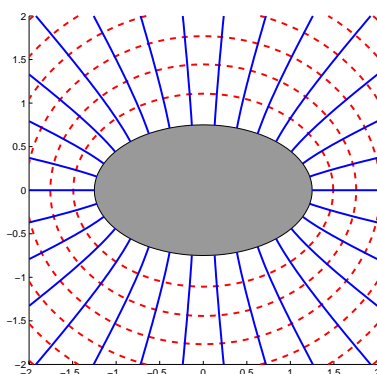


Figure 1: Initial mapping as given by $g(\omega, 0) = \omega + c/\omega$, with $c = 0.25$.
Image courtesy of Jaehyuk Choi. Used with permission.

2.3 M-Fold Perturbation of the Circle

We consider now the following mapping:

$$g(\omega, t) = A_1(t)\omega + \frac{A_{1-M}(t)}{\omega^{M-1}}$$

Again, without loss of generality, we will write:

$$\begin{aligned} A_1(0) &= 1 \\ A_{1-M}(0) &= \frac{c}{M-1} \end{aligned}$$

Let us see what the following mapping gives:

$$g(\omega, 0) = \omega + \frac{c}{(M-1)\omega^{M-1}} \quad (6)$$

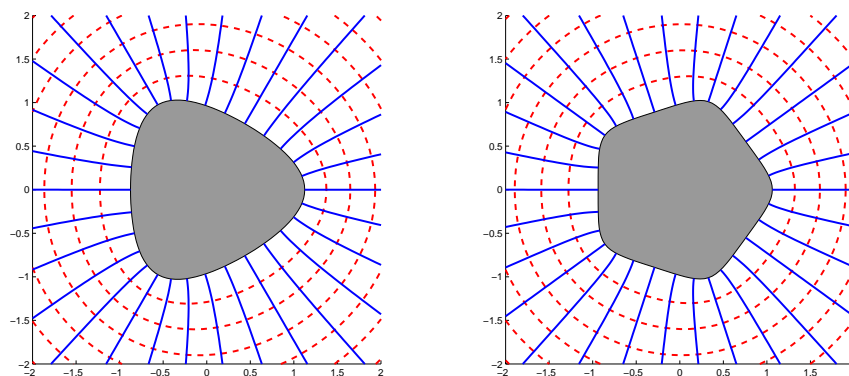


Figure 2: Mapping consisting of an M-fold perturbation of a circle, as given by (6): on the left, $M = 3$, on the right $M = 5$. For both, $c = 0.25$. Image courtesy of Jaehyuk Choi. Used with permission.

We evaluate:

$$g'(\omega, 0) = 1 - c\omega^{-M}$$

indicating that if $c \geq 1$, the mapping loses its conformality in the domain $|\omega| \geq 1$ at M points on a circle: $\omega_n = c^{1/M} e^{i2\pi n/M}$, $n = 0, 1, \dots, M-1$. However, if $c < 1$, conformality is never lost.

The area equation (4) gives:

$$|A_1(t)|^2 - (M-1)|A_{1-M}(t)|^2 = 2t + 1 - \frac{c^2}{M-1}$$

and (5) gives:

$$-(M-1)\dot{A}_1(t)A_{1-M}(t) + A_1(t)\dot{A}_{1-M}(t) = 0 \quad \Rightarrow \quad (M-1)\frac{\dot{A}_1(t)}{A_1(t)} = \frac{\dot{A}_{1-M}(t)}{A_{1-M}(t)}$$

which gives after integration:

$$A_1(t)^{M-1} = \frac{M-1}{c} A_{1-M}(t)$$

We finally get:

$$g(\omega, t) = A_1(t) \left[\omega + \frac{cA_1(t)^{M-2}}{(M-1)\omega^{M-1}} \right]$$

where $A_1(t)$ solves:

$$|A_1(t)|^2 - \frac{c^2}{M-1}|A_1(t)|^{2M-2} = 2t + 1 - \frac{c^2}{M-1}$$

Following the previous discussion, we want to check when the mapping loses its conformality on $|\omega| = 1$, that is we want to find t_c such that $cA_1(t_c)^{M-2} = 1$. Plugging this relation into the equation above gives:

$$t_c = \frac{1}{2} \left[\frac{c^2}{M-1} + \frac{M-2}{M-1} \left(\frac{1}{c} \right)^{\frac{2}{M-2}} - 1 \right]$$

which has a positive finite value for any $M \geq 3$ and $0 < c < 1$. Thus, the mapping will lose its conformality after a finite time, as illustrated in figure 3.

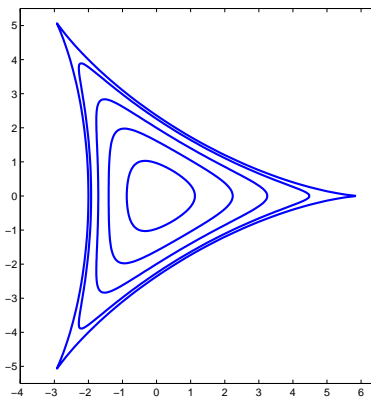


Figure 3: Evolution of a mapping consisting of a 3-fold perturbation of a circle ($M = 3$). In this case, $c = 0.25$ so that $t_c = 3.51563$. The contours are plotted at the times $t = 0, 1, 2, 3$ and 3.51563 . Image courtesy of Jaehyuk Choi. Used with permission.

2.4 Additional Remark

For a mapping with a finite Laurent expansion:

$$g(\omega, t) = \sum_{n=-N}^1 A_n(t) \omega^n \quad \text{with } N < \infty$$

we can still show that (5) leads to:

$$A_1(t)^N = C^{\text{te}} A_{-N}(t) \tag{7}$$

On the other hand, the area rule (4) gives:

$$|A_1(t)|^2 - \sum_{n=1}^N n |A_{-n}(t)|^2 = 2t + \text{const}$$

showing that $A_1(t)$, which gives the only positive term in the left hand side term, must grow at least as \sqrt{t} . Then from (7), we conclude that $A_{-N}(t)$ grows at least as $t^{N/2}$ which is in contradiction with the area rule. In order to solve this, the mapping must reach a singularity in a finite time t_c , as in the previous section.