# 2.003SC Recitation 11 Notes: Double Pendulum System

## **EIGENVALUES and EIGENVECTORS**

Consider the following matrix equation,

$$\underline{Ax} = \lambda \underline{x} \tag{1}$$

where <u>A</u> is a matrix of size nxn, <u>x</u> is a vector of length n, and  $\lambda$  is a scalar

For a given matrix, <u>A</u>, the values of  $\lambda_i$  and  $\underline{x}_i$ , i = 1, ..., n that satisfy the above equation are called (the matrix's) **eigenvalues** and **eigenvectors**, respectively.

Eigenvalues and eigenvectors are a very important and valuable concept that arises in many technical fields, especially vibrations.

Consequently, well-established, robust computational procedures exist for evaluating the eigenvalues and eigenvectors of a matrix.

#### **Connection to Vibrations**

Recall the matrix form of the equations of motion for an n-degree-of-freedom system,

$$\underline{M\ddot{x}} + \underline{Kx} = 0$$

This can be re-written as

$$\underline{\ddot{x}} + \underline{M}^{-1}\underline{K}\underline{x} = 0$$

or

$$\underline{M}^{-1}\underline{K}\underline{x} = -\underline{\ddot{x}}$$

Recall that, for harmonic motion

$$\underline{\ddot{x}} = -\omega^2 \underline{x}$$

So, the matrix equation has the same form as (1) above, i.e. can be seen to be an eigenvalue problem.

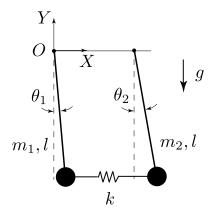
$$(\underline{M}^{-1}\underline{K})\underline{x} = (\omega^2)\underline{x}$$

where

- $\underline{A} = \underline{M}^{-1}\underline{K}$  is the system matrix
- the eigenvalues,  $\lambda_i$  , are the natural frequencies,  ${\omega_i}^2$
- the eigenvectors,  $\underline{x}_i$  , are the natural modes

#### Double Pendulum System - Problem Statement

Consider a system of two masses and one spring as shown in the figure below. Note that  $\theta_1$  and  $\theta_2$  are small-angle displacements.



The system's equations of motion are

$$\ddot{\theta}_1 + \left(\frac{g}{l} + \frac{k}{m_1}\right)\theta_1 - \left(\frac{k}{m_1}\right)\theta_2 = 0$$
$$\ddot{\theta}_2 - \left(\frac{k}{m_2}\right)\theta_1 + \left(\frac{g}{l} + \frac{k}{m_2}\right)\theta_2 = 0$$

For the special case where  $m_1 = m_2 = m$ ,

- Write the equations of motion in matrix notation.
- Find the characteristic equation
- Find the natural frequencies and natural modes

### Double Pendulum System - Solution

### EQUATIONS OF MOTION IN MATRIX NOTATION

$$\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \ddot{\theta}_1 \\ \ddot{\theta}_2 \end{bmatrix} + \begin{bmatrix} \left(\frac{g}{l} + \frac{k}{m_1}\right) & -\frac{k}{m_1} \\ -\frac{k}{m_2} & \left(\frac{g}{l} + \frac{k}{m_2}\right) \end{bmatrix} \begin{bmatrix} \theta_1 \\ \theta_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
(1)

or

$$\underline{M\ddot{x}} + \underline{Kx} = 0$$

Setting  $m_1 = m_2 = m$ , the equations of motion are

$$\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \ddot{\theta}_1 \\ \ddot{\theta}_2 \end{bmatrix} + \begin{bmatrix} (\frac{g}{l} + \frac{k}{m}) & -\frac{k}{m} \\ -\frac{k}{m} & (\frac{g}{l} + \frac{k}{m}) \end{bmatrix} \begin{bmatrix} \theta_1 \\ \theta_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
(2)

### CHARACTERISTIC EQUATION

Assume the two masses undergo <u>harmonic motion</u>, i.e. they oscillate with the <u>same frequency</u>,  $\omega$ , albeit different amplitudes,  $a_1, a_2$ .

$$\begin{bmatrix} \theta_1 \\ \theta_2 \end{bmatrix} = \begin{bmatrix} a_1 \\ a_2 \end{bmatrix} \cos(\omega t - \phi) \tag{3}$$

$$\begin{bmatrix} \ddot{\theta}_1\\ \ddot{\theta}_2 \end{bmatrix} = -\omega^2 \begin{bmatrix} a_1\\ a_2 \end{bmatrix} \cos(\omega t - \phi)$$
(4)

Substituting (3) and (4) into (2), we obtain,

$$\begin{bmatrix} -\omega^2 & 0\\ 0 & -\omega^2 \end{bmatrix} \begin{bmatrix} a_1\\ a_2 \end{bmatrix} \cos(\omega t - \phi) + \begin{bmatrix} \left(\frac{g}{l} + \frac{k}{m}\right) & -\frac{k}{m}\\ -\frac{k}{m} & \left(\frac{g}{l} + \frac{k}{m}\right) \end{bmatrix} \begin{bmatrix} a_1\\ a_2 \end{bmatrix} \cos(\omega t - \phi) = \begin{bmatrix} 0\\ 0 \end{bmatrix}$$

Dividing by  $cos(\omega t - \phi)$ , we obtain

$$\begin{bmatrix} -\omega^2 + \left(\frac{g}{l} + \frac{k}{m}\right) & -\frac{k}{m} \\ -\frac{k}{m} & -\omega^2 + \left(\frac{g}{l} + \frac{k}{m}\right) \end{bmatrix} \begin{bmatrix} a_1 \\ a_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
(5)

Setting the determinant equal to zero produces the CHARACTERISTIC EQUATION/POLYNOMIAL.

$$h^2 - 2h\omega^2 + \omega^4 - \left(\frac{k}{m}\right)^2 = 0$$
 where  $h = \frac{g}{l} + \frac{k}{m}$ 

### NATURAL FREQUENCIES AND NATURAL MODES

Applying the quadratic formula to the characteristic equation,

$$\omega^2 = h \pm \frac{k}{m} = \left(\frac{g}{l} + \frac{k}{m}\right) \pm \frac{k}{m}$$

or

$$\omega_1 = \sqrt{\frac{g}{l}}$$
;  $\omega_2 = \sqrt{\frac{g}{l} + 2\frac{k}{m}}$ 

From the first row of (5),

$$\left(-\omega^2 + \frac{g}{l} + \frac{k}{m}\right)a_1 - \frac{k}{m}a_2 = 0$$

we can obtain the formula for the <u>natural modes</u>,

$$\frac{a_2}{a_1} = \frac{-m\omega^2 + 2k}{k}$$

which we evaluate at each of the natural frequencies,

$$\omega_1 = \sqrt{\frac{g}{l}} \qquad \longrightarrow \qquad \frac{a_2}{a_1} = 1$$

$$\omega_2 = \sqrt{\frac{g}{l} + 2\frac{k}{m}} \qquad \rightarrow \qquad \frac{a_2}{a_1} = -1$$

#### **General Solution**

In general, however, (i.e. for <u>arbitrary</u> initial conditions), the system's free response will contain  $\underline{BOTH}$  natural frequencies,

$$\begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = A_1 \begin{bmatrix} 1 \\ 1 \end{bmatrix} \cos(\omega_1 t - \phi_1) + A_2 \begin{bmatrix} 1 \\ -1 \end{bmatrix} \cos(\omega_2 t - \phi_2)$$

where  $A_1, A_2, \phi_1, \phi_2$  are determined by initial conditions.

This general response can appear to be very irregular, with little discernible pattern. When the natural frequencies are close together, "beating" behavior can be observed.

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