### 2.003SC Recitation 11 Notes: Double Pendulum System

## EIGENVALUES and EIGENVECTORS

Consider the following matrix equation,

$$
\begin{equation*}
\underline{A x}=\lambda \underline{x} \tag{1}
\end{equation*}
$$

where $\underline{A}$ is a matrix of size $n x n, \underline{x}$ is a vector of length $n$, and $\lambda$ is a scalar
For a given matrix, $\underline{A}$, the values of $\lambda_{i}$ and $\underline{x}_{i}, i=1, \ldots, n$ that satisfy the above equation are called (the matrix's) eigenvalues and eigenvectors, respectively.

Eigenvalues and eigenvectors are a very important and valuable concept that arises in many technical fields, especially vibrations.
Consequently, well-established, robust computational procedures exist for evaluating the eigenvalues and eigenvectors of a matrix.

## Connection to Vibrations

Recall the matrix form of the equations of motion for an n-degree-of-freedom system,

$$
\underline{M \ddot{x}}+\underline{K x}=0
$$

This can be re-written as

$$
\underline{\ddot{x}}+\underline{M}^{-1} \underline{K x}=0
$$

or

$$
\underline{M}^{-1} \underline{K x}=-\underline{\ddot{x}}
$$

Recall that, for harmonic motion

$$
\underline{\ddot{x}}=-\omega^{2} \underline{x}
$$

So, the matrix equation has the same form as (1) above, i.e. can be seen to be an eigenvalue problem.

$$
\left(\underline{M}^{-1} \underline{K}\right) \underline{x}=\left(\omega^{2}\right) \underline{x}
$$

where

- $\underline{A}=\underline{M}^{-1} \underline{K}$ is the system matrix
- the eigenvalues, $\lambda_{i}$, are the natural frequencies, $\omega_{i}{ }^{2}$
- the eigenvectors, $\underline{x}_{i}$, are the natural modes


## Double Pendulum System - Problem Statement

Consider a system of two masses and one spring as shown in the figure below. Note that $\theta_{1}$ and $\theta_{2}$ are small-angle displacements.


The system's equations of motion are

$$
\begin{aligned}
& \ddot{\theta}_{1}+\left(\frac{g}{l}+\frac{k}{m_{1}}\right) \theta_{1}-\left(\frac{k}{m_{1}}\right) \theta_{2}=0 \\
& \ddot{\theta}_{2}-\left(\frac{k}{m_{2}}\right) \theta_{1}+\left(\frac{g}{l}+\frac{k}{m_{2}}\right) \theta_{2}=0
\end{aligned}
$$

$\underline{\text { For the special case where } m_{1}=m_{2}=m, ~}$

- Write the equations of motion in matrix notation.
- Find the characteristic equation
- Find the natural frequencies and natural modes


## Double Pendulum System - Solution

## EQUATIONS OF MOTION IN MATRIX NOTATION

$$
\left[\begin{array}{ll}
1 & 0  \tag{1}\\
0 & 1
\end{array}\right]\left[\begin{array}{l}
\ddot{\theta}_{1} \\
\ddot{\theta}_{2}
\end{array}\right]+\left[\begin{array}{cc}
\left(\frac{g}{l}+\frac{k}{m_{1}}\right) & -\frac{k}{m_{1}} \\
-\frac{k}{m_{2}} & \left(\frac{g}{l}+\frac{k}{m_{2}}\right)
\end{array}\right]\left[\begin{array}{l}
\theta_{1} \\
\theta_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
$$

or

$$
\underline{M \ddot{x}}+\underline{K x}=0
$$

Setting $m_{1}=m_{2}=m$, the equations of motion are

$$
\left[\begin{array}{ll}
1 & 0  \tag{2}\\
0 & 1
\end{array}\right]\left[\begin{array}{l}
\ddot{\theta}_{1} \\
\ddot{\theta}_{2}
\end{array}\right]+\left[\begin{array}{cc}
\left(\frac{g}{l}+\frac{k}{m}\right) & -\frac{k}{m} \\
-\frac{k}{m} & \left(\frac{g}{l}+\frac{k}{m}\right)
\end{array}\right]\left[\begin{array}{l}
\theta_{1} \\
\theta_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
$$

## CHARACTERISTIC EQUATION

Assume the two masses undergo harmonic motion, i.e. they oscillate with the same frequency, $\omega$, albeit different amplitudes, $a_{1}, a_{2}$.

$$
\begin{gather*}
{\left[\begin{array}{l}
\theta_{1} \\
\theta_{2}
\end{array}\right]=\left[\begin{array}{l}
a_{1} \\
a_{2}
\end{array}\right] \cos (\omega t-\phi)}  \tag{3}\\
{\left[\begin{array}{l}
\ddot{\theta}_{1} \\
\ddot{\theta}_{2}
\end{array}\right]=-\omega^{2}\left[\begin{array}{l}
a_{1} \\
a_{2}
\end{array}\right] \cos (\omega t-\phi)} \tag{4}
\end{gather*}
$$

Substituting (3) and (4) into (2), we obtain,

$$
\left[\begin{array}{cc}
-\omega^{2} & 0 \\
0 & -\omega^{2}
\end{array}\right]\left[\begin{array}{l}
a_{1} \\
a_{2}
\end{array}\right] \cos (\omega t-\phi)+\left[\begin{array}{cc}
\left(\frac{g}{l}+\frac{k}{m}\right) & -\frac{k}{m} \\
-\frac{k}{m} & \left(\frac{g}{l}+\frac{k}{m}\right)
\end{array}\right]\left[\begin{array}{l}
a_{1} \\
a_{2}
\end{array}\right] \cos (\omega t-\phi)=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
$$

Dividing by $\cos (\omega t-\phi)$, we obtain

$$
\left[\begin{array}{cc}
-\omega^{2}+\left(\frac{g}{l}+\frac{k}{m}\right) & -\frac{k}{m}  \tag{5}\\
-\frac{k}{m} & -\omega^{2}+\left(\frac{g}{l}+\frac{k}{m}\right)
\end{array}\right]\left[\begin{array}{l}
a_{1} \\
a_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]
$$

Setting the determinant equal to zero produces the CHARACTERISTIC EQUATION/POLYNOMIAL.

$$
h^{2}-2 h \omega^{2}+\omega^{4}-\left(\frac{k}{m}\right)^{2}=0 \quad \text { where } \quad h=\frac{g}{l}+\frac{k}{m}
$$

## NATURAL FREQUENCIES AND NATURAL MODES

Applying the quadratic formula to the characteristic equation,

$$
\omega^{2}=h \pm \frac{k}{m}=\left(\frac{g}{l}+\frac{k}{m}\right) \pm \frac{k}{m}
$$

or

$$
\omega_{1}=\sqrt{\frac{g}{l}} \quad ; \quad \omega_{2}=\sqrt{\frac{g}{l}+2 \frac{k}{m}}
$$

From the first row of (5),

$$
\left(-\omega^{2}+\frac{g}{l}+\frac{k}{m}\right) a_{1}-\frac{k}{m} a_{2}=0
$$

we can obtain the formula for the natural modes,

$$
\frac{a_{2}}{a_{1}}=\frac{-m \omega^{2}+2 k}{k}
$$

which we evaluate at each of the natural frequencies,

$$
\begin{array}{ccc}
\omega_{1}=\sqrt{\frac{g}{l}} & \rightarrow & \frac{a_{2}}{a_{1}}=1 \\
\omega_{2}=\sqrt{\frac{g}{l}+2 \frac{k}{m}} & \rightarrow & \frac{a_{2}}{a_{1}}=-1
\end{array}
$$

## General Solution

In general, however, (i.e. for arbitrary initial conditions), the system's free response will contain BOTH natural frequencies,

$$
\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right]=A_{1}\left[\begin{array}{l}
1 \\
1
\end{array}\right] \cos \left(\omega_{1} t-\phi_{1}\right)+A_{2}\left[\begin{array}{c}
1 \\
-1
\end{array}\right] \cos \left(\omega_{2} t-\phi_{2}\right)
$$

where $A_{1}, A_{2}, \phi_{1}, \phi_{2}$ are determined by initial conditions.

This general response can appear to be very irregular, with little discernible pattern. When the natural frequencies are close together, "beating" behavior can be observed.

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