

## Hints to Assignment #8 -- 8.022

### [1]/[2]/[4] Fields from moving particles: Purcell 5.13/5.14/5.19

These three problems deal with charged particles that are in motion and their associated electric fields. Please review sections 5.6 and 5.7 (p. 181-189) of Purcell's chapter 5 BEFORE you attempt these problems. Appendix B (p.459) of Purcell's gives a more thorough picture of the phenomenon.

- 5.13: The shown field lines can be split into two regions: one corresponding to a particle at rest and converging at  $x=0$  (how can you tell? watch the angles between the field lines- are they equal?) and one corresponding to a particle that moves and converging at  $x=12\text{cm}$ . Assume that any change of its velocity occurs within a tiny  $dt=0$ .
- 5.14: The colliding  $e^+ e^-$  meet and stop at  $x=0$  and thus a null field should emanate from them starting at  $t=0$ . This information though, i.e., that the field has vanished travels out at  $c$ , so at any instance  $t>0$  ONLY a sphere of radius  $ct$  will reflect this event. The region outside that sphere will still "see" the field resulting from the two electrons in the positions they WOULD have been at time  $t$ , i.e., the  $e^+$  at  $+vt$  and the  $e^-$  at  $-vt$ . Thus outside the causal sphere, lines should emerge pointing back to these two locations. Draw these lines at two instances  $t_1$  and  $t_2>t_1$ . Notice that the "width" of the lines wrapping up the causal sphere remains the same as it reflects the time the two electrons spent to dispose their kinetic energy and thus come at rest.
- 5.19: As the proton modifies its travel plan at  $t=0$  by colliding with the nucleus (and inverting) each velocity direction, at time  $T$  only a causal sphere of radius  $R=cT$  "knows" about this change. Thus the world outside this sphere, knows that the proton travels at  $u=-c/2$  and "expects" the proton to be at  $X=-uT$  at time  $T$  ... while the world inside this sphere got the news already that the proton flipped velocity sign and at time  $T$  should be at  $X=+uT$ . Draw the causal sphere  $R$  and the field lines converging at  $X=-uT$  and  $X=+uT$ . Lines should "break" right on the sphere's surface (or circle's periphery, given that 2-dim pictures are always easier to draw).

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### [3] Line charges and currents, Purcell 5.18

As the problem suggests, section 5.9 has all the steps needed in deriving the transformations. Start with  $I'=\lambda' \beta' c$  and expressing the primed quantities in terms of the unprimed.

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The next two problems require the understanding of the relationship between currents and magnetic fields as they are expressed by the Biot-Savart and Ampere's laws. This is discussed in the assigned reading of chapter 6 of your textbook. We will be talking on these in our Tuesday, Oct 29 class meeting. Here are some highlights of the Biot-Savart and Ampere's laws.

- You may view Biot-Savart law as the corresponding to Coulomb's law ( $E \sim 1/r^2$ ) that relates the source (currents=moving charges) and result (magnetic field). Biot-Savart wrote down what is the infinitesimal magnetic field (vector!!)  $d\mathbf{B}$  due to an infinitesimal part  $d\mathbf{L}$  (vector) of a current carrier  $I$  (scalar) at distance  $r$  (vector) from it.

This states:  $d\mathbf{B} = I d\mathbf{L} \times \mathbf{r} / (cr^3)$

- The key feature of Biot-Savart's Law is that the resulting magnetic field is given by the CROSS PRODUCT of two vectors,  $d\mathbf{L}$  and  $\mathbf{r}$ .
- The above law can be used a priori to derive the magnetic field (magnitude AND direction) at a point in space  $\mathbf{r}$  (vector) and for ANY current-carrying distribution by performing a mere integration over the current carrier.
- A feature of the above derived field is given by Ampere's Law which summarizes a concept very close to Gauss's one: the integral of  $\mathbf{B}$  (vector) around a closed loop is proportional to the enclosed current by the loop. In the CGS system, the constant of proportionality is given by  $4\pi/c$  (in MKSA it is  $\mu_0$ ).
- Applying Ampere's Law for calculating  $\mathbf{B}$  is subject to the same limitation of Gauss's law for calculating  $\mathbf{E}$ , i.e., only highly symmetric distributions of currents may allow the derivation of  $\mathbf{B}$  using it. One should come up with a loop over which  $\mathbf{B}$  has either a constant value and constant direction either in whole or in parts so that the  $\mathbf{B} \cdot d\mathbf{L}$  (inner product of vectors) can be written in terms of magnitudes. What Ampere's law dictates is that the sum of all the  $\mathbf{B} \cdot d\mathbf{L}$ 's along a path is proportional to the enclosed currents. In order to establish the enclosed currents you should imagine ANY FINITE surface that TERMINATES on the loop. In most cases, the simpler the finite surface the better for you to easily draw it in paper and envision what is going on. HATCH in a distinct style that surface and identify its intersection with ANY CURRENT distribution. You now have to sum up the CURRENT that overlaps with the hatched area. Current is a scalar quantity, i.e. a signed number. FIRST you have to assign a direction along which you will circle around the amperian loop. THIS IS ARBITRARY. Remember that  $\mathbf{B} \cdot d\mathbf{L} = -\mathbf{B} \cdot d\mathbf{L}'$  where  $d\mathbf{L}$  and  $d\mathbf{L}'$  have opposite directions. You will then have to assign the "+" to the currents that flow along the direction of your right hand thumb when your (right hand) palm traces the previously assigned direction of your amperian loop. Any currents flowing in the opposite direction get the "-" sign. You are now ready to build the (scalar) quantity of the RHS of Ampere's equation, i.e., the total current enclosed by the loop. WATCH out for cases that the current is described by DISTRIBUTION (i.e., current densities) that are NOT uniform- in this case a surface integral has to be performed.