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8.512 Theory of Solids II
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8.512 Theory of Solids

Problem Set 12

Due May 13, 2009

Consider a Fermi gas with dispersion ϵ_k and a repulsive interaction $U\delta(\mathbf{r})$. Now if $N(0)U > 1$, we find in mean field theory the spontaneous appearance of the order parameter:

$$\Delta = U\langle n_\uparrow - n_\downarrow \rangle ,$$

and the splitting of the up- and down-spin bands, so that the mean field Hamiltonian is

$$H_{\text{MF}} = \sum_k \left(\tilde{\epsilon}_{k\uparrow} c_{k\uparrow}^\dagger c_{k\uparrow} + \tilde{\epsilon}_{k\downarrow} c_{k\downarrow}^\dagger c_{k\downarrow} \right) ,$$

where

$$\begin{aligned} \tilde{\epsilon}_{k\uparrow} &= \epsilon_k - \Delta/2 \\ \tilde{\epsilon}_{k\downarrow} &= \epsilon_k + \Delta/2 . \end{aligned}$$

1. Consider the **transverse** spin susceptibility. Instead of $\chi_x = dM_x/dH_x$, where $\mathbf{M} = -2\mu_B(\boldsymbol{\sigma}/2)$ it is more convenient to consider the response to H_+ which couples to a spin flip excitation, i.e., $\chi_\perp = d(M_+/dH_+)$ where $M_+ = \frac{1}{2}(M_x + iM_y)$ and $H_+ = H_x + iH_y$. Show that the response function for the mean field Hamiltonian is given by $\chi_\perp^0 = \mu_B^2 \Gamma_0$ where

$$\Gamma_0(q, \omega) = \sum_k \frac{f(\tilde{\epsilon}_{k+q,\downarrow}) - f(\tilde{\epsilon}_{k,\uparrow})}{\omega - \tilde{\epsilon}_{k+q,\downarrow} + \tilde{\epsilon}_{k,\uparrow} + i\eta} .$$

This is the generalization of the Lindhard function to a spin split band.

2. Now include the interaction term in the response to the transverse field in a self consistent field approximation. Show that

$$\chi_{\perp}(q, \omega) = \frac{\mu_B^2 \Gamma_0(q, \omega)}{1 - U\Gamma_0(q, \omega)} .$$

3. The poles of the numerator in χ_{\perp} describe the single particle-hole excitations. Sketch the region in (ω, q) space where $Im\chi_{\perp} \neq 0$ due to these excitations.
4. The other pole in $\chi_{\perp}(q, \omega)$ occurs when the denominator vanishes. Calculate the dispersion of this pole which we identify as the spin wave excitation as follows:
 - (a) Show that at $q = \omega = 0$, the denominator vanishes. [Hint: the condition $1 - U\Gamma_0 = 0$ is the same as the self-consistency equation for $\Delta(T)$.]
 - (b) Expand $\Gamma_0(q, \omega)$ for small q, ω and show that the location of the pole of χ_{\perp} is given by $\omega(q) = Cq^2$. Note that unlike the Lindhard function for free fermions, the existence of the gap Δ makes the expansion well behaved.